Search for metastable heavy charged particles with large ionization energy loss in pp collisions at root s=13 TeV using the ATLAS experiment

Aaboud, M.; Aad, G.; Abbott, B.; Abdallah, J.; Abdinov, O.; Abelaos, B.; Aben, R.; AbouZeid, O.S.; Nitzan, Abraham; Abramowicz, H.; Dam, Mogens; Hansen, Jørn Dines; Hansen, Jørgen Beck; Xella, Stefania; Hansen, Peter Henrik; Petersen, Troels Christian; Alonso Diaz, Alejandro; Wiglesworth, Graig; Pingel, Almut Maria; Monk, James William; Thomsen, Lotte Ansgaard; Pedersen, Lars Egholm; Løvschall-Jensen, Ask Emil; Galster, Gorm Aske Gram Krohn

Published in: Physical Review D (Particles, Fields, Gravitation and Cosmology)

DOI: 10.1103/PhysRevD.93.112015

Publication date: 2016

Document version Publisher's PDF, also known as Version of record


Download date: 06. Oct. 2023
I. INTRODUCTION

Massive, long-lived particles (LLPs) are predicted by a wide range of physics models that extend the Standard Model (SM). LLPs arise in proposed solutions to the gauge hierarchy problem [1], including supersymmetric (SUSY) models that either violate [2–4] or conserve [5–12] R-parity. The lifetime of these particles depends on the mass difference between the particle and the lightest stable SUSY particle, and on the size of any R-parity-violating coupling [13].

Because of their large mass, LLPs are expected to be slow (β significantly below 1) and, if charged, to have a specific ionization higher than any SM particles of unit charge at high momenta. The Pixel subsystem [14] of the ATLAS detector [15] provides measurements of ionization energy loss (dE/dx) for charged particles and can be used to distinguish such highly ionizing particles from SM particles. In this paper, this information is used to search for LLPs using data collected at √s = 13 TeV, extending in reach a study performed on data collected at √s = 8 TeV [16]. The increase in probability for high-energy parton-parton collisions significantly enhances the production cross section of LLPs via the strong interaction.

Together with this enhancement, the search sensitivity has been considerably extended thanks to the upgrades to the Pixel detector (see Sec. II) and improvements to the analysis (see Sec. VI).

Many strategies have been explored to find LLPs at the LHC. A search using the complete √s = 8 TeV data set recorded by the ATLAS detector [17] combined Pixel dE/dx information with time-of-flight measurements, made in both the calorimeters and the muon spectrometer, to search for LLPs that do not decay within the detector. Charginos that are nearly mass degenerate with the lightest stable SUSY particle have been searched for in ATLAS looking for LLPs with short lifetimes by looking for secondary vertices [21,22] and for disappearing tracks [23].

The current study addresses many different models of new physics, in particular those that predict the production of metastable massive particles with O(ns) lifetime at LHC energies, such as mini-split SUSY [11,24,25] or anomaly-mediated supersymmetry-breaking models [26,27]. A metastable gluino with a mass of approximately 1 TeV would be compatible with the measured Higgs boson mass according to mini-split SUSY models, which also predict squarks with masses of 10^3–10^5 TeV, therefore making the...
gluino the only observable sparticle produced through the strong interaction at LHC energies. Results are presented assuming the production of $R$-hadrons as composite colorless states of a gluino together with SM quarks or gluons and the subsequent decay, on nanosecond timescales, of the gluino to a stable neutralino and a $q\bar{q}$ pair [28]. This search could, in principle, also be sensitive to the production of LLPs through weak interactions, but these models are not considered here as the cross sections are much lower.

In the rest of this paper, $R$-hadrons with lifetimes consistent with decaying inside the ATLAS detector are referred to as metastable; the search strategy focuses on decays occurring within the active detector volume, which covers lifetimes from around 1 ns to several tens of ns. $R$-hadrons whose lifetime disfavors decay within the ATLAS detector are termed stable. The approach described in this paper allows direct observation of charged $R$-hadrons which are either stable or metastable and traverse the first seven layers of the ATLAS inner tracking system.

II. ATLAS DETECTOR AND PIXEL $dE/dx$ MEASUREMENT

The ATLAS detector\(^1\) consists of a tracker, for measuring the trajectories of charged particles, surrounded by a solenoid magnet, followed by calorimeters for measuring the energy of particles that interact electromagnetically or hadronically. A muon spectrometer immersed in a toroidal magnetic field surrounds the calorimeters, and provides tracking for muons. The detector is hermetic and can therefore measure the missing transverse momentum associated with each event. A complete description of the ATLAS detector can be found in Ref. [15]. The tracker is made of three detector systems. Moving from the solenoid magnet toward the beam collision region, there is first a transition radiation tracker with approximately 400,000 channels [29], followed by a silicon microstrip detector (SCT) with about 6 million channels [30], and at the innermost radius sits a roughly 92 million channel silicon Pixel detector, which is crucial for this measurement, and is described below in some detail.

The current Pixel detector provides at least four precision measurements for each track in the region $|\eta| < 2.5$ at radial distances of 3.4 to 13 cm from the LHC beam line. After Run 1, the ATLAS Pixel detector was upgraded with the insertion of an additional layer, the Insertable B-Layer (IBL), which was installed inside the Run 1 Pixel detector, mounted on a new beam pipe of smaller diameter. The IBL

\(^1\)ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point in the center of the detector and the $z$ axis coinciding with the axis of the beam pipe. The $x$ axis points from the interaction point to the center of the LHC ring, and the $y$ axis points upward. Cylindrical coordinates $(r, \phi)$ are used in the transverse plane, $\phi$ being the azimuthal angle around the $z$ axis. The pseudorapidity is defined in terms of the polar angle $\theta$ as $\eta = -\ln \tan(\theta/2)$. 

has smaller area pixels, reduced thickness, and different coding electronics, which provides charge measurements with lower resolution and dynamic range than the other Pixel layers. It therefore requires a separate treatment of the $dE/dx$ information. At normal incidence, the average charge released by a minimum ionizing particle (MIP) in a pixel sensor is $\approx 20000 e^-$ ($\approx 16000 e^-$ for IBL) and the charge threshold is set to $3500 \pm 40 e^-$ ($2500 \pm 40 e^-$ for IBL). Signals are accepted if they are larger than this threshold. They are then matched to a specific beam crossing. The hit efficiency under these conditions exceeds 99%. When detector data are read out, the time over threshold (ToT), i.e. the time interval with the signal above the threshold, is digitized with 8 bits (4 bits for IBL). The ToT is approximately proportional to the ionization charge [31] and its dynamic range corresponds to 8.5 times (1.5 times for IBL) the average charge released by a MIP for a track normal to the silicon detectors which deposits all its ionization charge in a single pixel. If this value is exceeded, the hit charge information is either underestimated in the IBL, where the electronics signals the excess with an overflow bit, or lost in the other pixel layers.

The charge released by a track crossing the Pixel detector is rarely contained within just one pixel; neighboring pixels registering hits are joined together using a connected component analysis [32,33] to form clusters. The charge of a cluster is calculated by summing the charges of all

![FIG. 1. Distribution of the $dE/dx$ computed with the Pixel algorithm in data (black dots) and Pixel without IBL (red triangles) for minimum bias events. Tracks were selected as in Ref. [34], with the additional requirements of $p_T > 3$ GeV and at least two associated pixel clusters. The corresponding $dE/dx$ distributions for simulation with and without the IBL are also plotted, respectively, as dashed and dotted lines. The simulation overestimates the $dE/dx$, which is taken into account in the systematic uncertainty evaluation. The addition of IBL clusters reduces the tails at both low and high $dE/dx$. The fraction of tracks with $dE/dx < 0.5 (\geq 1.8)$ MeV g$^{-1}$ cm$^2$ decreases with the use of the IBL charge information from 0.065% to 0.018% (from 0.89% to 0.44%).]
pixels belonging to the cluster after calibration corrections; the IBL overflow information is accounted for separately. The dE/dx is estimated using the average of the individual cluster ionization measurements (charge collected in the cluster per unit track length in the sensor), for the clusters associated with a track. To reduce the effect of the tails of the Landau distribution, the average is evaluated after removing the highest dE/dx cluster, or the two highest dE/dx clusters in the rare case of more than four clusters on the track. A track is considered to have a good ionization measurement only if there are at least two remaining clusters. Including the IBL charge information reduces the tails of the dE/dx distribution, as illustrated in Fig. 1, and increases the fraction of tracks for which the ionization measurement is valid from 77% to 91%. The most probable value of the measured dE/dx is 1.12 MeV g\(^{-1}\) cm\(^{-2}\) for a minimum ionizing particle. This value is obtained with a Gaussian fit to the data shown in Fig. 1, in a ±1 standard deviation (0.13 MeV g\(^{-1}\) cm\(^{-2}\) ) range around the peak. The βγ measurable with the current dE/dx method cannot be smaller than approximately 0.3 for particles with unit charge because of the ToT dynamic range.

III. MASS CALCULATION

The average energy loss of massive, charged particles in matter is expected to follow the Bethe-Bloch distribution, which can be expressed as a function of the βγ of the LLPs. Their mass can be derived from a fit of the measured specific energy loss and the reconstructed momentum to a parametric Bethe-Bloch distribution in the range 0.3 < βγ < 1.5. This range overlaps the expected average βγ of LLPs produced at the LHC, which decreases with the particle mass from ⟨βγ⟩ ≈ 2.0 at 100 GeV to ⟨βγ⟩ ≈ 0.5 at 1600 GeV. The parametric function describing the relationship between the most probable value of the energy loss (dE/dx)\(_{MPV}\) and βγ is

\[
(dE/dx)_{MPV}(βγ) = \frac{p_1}{β^2γ} \ln(1 + [p_2βγ]^{p_3}) - p_4. \tag{1}
\]

The \(p_i\), with \(i = 1...5\), calibration constants were measured in this data set using the peak values of three Crystal Ball functions, fitted to the dE/dx distribution observed in each of several low-momentum slices to model the energy loss distribution for pions, kaons and protons respectively, as described in Ref. [35].

The distribution of the dE/dx versus momentum is shown in Fig. 2 together with the fits to the pion, kaon and proton masses. The mass calculation performed through the dE/dx method is monitored by checking the stability of the proton mass measurement during the data taking.

A mass estimate \(M\) is obtained by numerically solving Eq. (1) for the unknown \(M\). In simulated R-hadron events, the reconstructed mass is found to reproduce well the generated mass up to masses around 1800 GeV, and a residual 3% correction is applied to the reconstructed mass to improve the agreement. For particles with higher masses (beyond the current sensitivity of this analysis), the reconstructed mass is observed to slightly underestimate the generated mass. The half width at half maximum of the reconstructed mass distribution increases with the mass value. This is due to the momentum measurement uncertainty dominating the mass resolution for masses greater than \(\approx 200\) GeV.

IV. R-HADRON SIMULATION

A number of samples of simulated signal events are used in this analysis to determine the expected LLP signal efficiencies and to estimate uncertainties in the efficiency. For stable R-hadrons, pair production of gluinos with masses between 800 GeV and 1800 GeV is simulated in \textsc{Pythia} 6.4.27 [36] with the \textsc{AueT2B} [37] set of tuned parameters for the underlying event and the CTEQ6L1 [38] parton distribution function (PDF) set, incorporating dedicated hadronization routines [39] to produce final states containing R-hadrons.

The cross sections are calculated at next-to-leading order (NLO) in the strong coupling constant including the resummation of soft-gluon emission at next-to-leading-logarithmic accuracy (NLO + NLL) [40–44] and assuming a squark mass of 10 TeV. The nominal cross-section predictions and the uncertainty are taken from an envelope of cross-section predictions using different PDF sets and factorization and renormalization scales, as described in Ref. [45].

Selected events containing R-hadrons undergo a full detector simulation [46], where interactions of R-hadrons with matter are handled by dedicated \textsc{Geant4} [47] routines based on the physics model described in Refs. [39,48].
model assumes, for each light valence quark, a cross section of 12 mb per nucleon and neglects the heavy parton, other than as a reservoir of kinetic energy. This hadronic scattering is described through a purely phase-space-driven approach. The probability for a gluino to form a gluon-gluino bound state is assumed to be 10% [5].

The simulation of samples of metastable gluino-based R-hadrons is performed similarly to the stable R-hadrons. The gluinos within R-hadrons are then required to decay via the process $\tilde{g} \rightarrow q\tilde{q}$, using PYTHIA 6. Gluino masses (m) between 400 and 3000 GeV are simulated, with the neutralino mass fixed to 100 GeV. This benchmark decay mode is chosen as it has the highest sensitivity among those studied in Ref. [16]. Heavier neutralino masses would give lower signal efficiencies, mainly because of the smaller missing transverse momentum available to trigger. The gluino lifetime (τ) is varied from 0.4 to 50 ns.

All simulated samples include a modeling of pileup, adding the expected number of minimum-bias proton-proton interactions from the same (in-time pileup) and nearby (out-of-time pileup) bunch crossings. Simulated events are reconstructed using the software used for the collision data.

To get a more accurate description of radiative effects than PYTHIA 6 provides, additional samples of gluinos are generated for all mass points using MG5_aMC@NLO [49], interfaced to the PYTHIA 8.186 [50] parton shower model, with the A14 [51] set of tuning parameters together with the NNPDF2.3LO [52] PDF set. The distribution of the transverse momentum of the gluino-gluino system simulated with PYTHIA 6 is then reweighted in such a way as to match that obtained in samples simulated with MG5_aMC@NLO. This is especially important for stable R-hadrons as detailed in Sec. VI.

V. SEARCH STRATEGY

This search is based on the signature of R-hadron events containing at least one reconstructed track with large measured $dE/dx$ and high transverse momentum ($p_T$). This signature optimizes the search for metastable R-hadrons with lifetimes of $O(\text{ns})$, as the selection requires that the R-hadron only live long enough to traverse the seven layers of the silicon detectors, corresponding to a distance from the beam axis of 37 cm. The search is also sensitive to R-hadrons with longer lifetimes, with a small dependence on the modeling of the interaction of R-hadrons in the calorimeter and muon systems.

The expected production and interaction of R-hadrons in the detector is described in detail in Ref. [53]. In the models studied in this search and described in Sec. IV, pair-produced gluinos hadronize into two colorless R-hadrons with the following charge states: approximately 33% events with two neutral R-hadrons, 47% events with one neutral and one charged R-hadron, and 20% events with two charged R-hadrons. The fragmentation is extremely hard and the R-hadrons are therefore isolated. As massive particles with $\beta < 1$, the charged R-hadrons are expected to deposit more ionization energy in the detector than typical particles of the same momentum. This ionization signature is used both to select signal events, as described in Sec. VI, and to calculate the mass of selected candidates, as described in Sec. III. The estimated candidate mass is used as the final search discriminant.

If the R-hadrons decay before or in the calorimeter, the decay products of R-hadrons in models studied here include quarks, which are identified as jets in the calorimeters (see Sec. VI), and a 100 GeV neutralino, which carries away significant momentum undetected. If R-hadrons traverse the calorimeter without decaying, they are expected to deposit significantly less energy in the calorimeters than SM hadrons of the same momentum, as the heavy parton carries almost all the momentum of the composite particle and limited kinetic energy is available for interactions between the light SM partons and the detector. In the models studied here, over 90% of stable R-hadrons, independent of mass, deposit less than 20 GeV of energy in the calorimeters. However, the electric charge of the R-hadron can change due to the hadronic interactions, even if little energy is lost, and the charge of the R-hadron leaving the calorimeter is expected to be decoupled from the charge at production [53]. As hadronic interactions are rare in the tracking detectors, the probability for an R-hadron to change its electric charge before the calorimeters is only around 3%.

Events are selected by a trigger requiring missing transverse momentum ($E_T^{\text{miss}}$), calculated from energy deposits in the calorimeter [54]. For events with R-hadrons that decay before or inside the calorimeter, undetected neutralinos contribute to the measured transverse momentum imbalance. As stable R-hadrons deposit little energy in the calorimeters, most stable R-hadron events are selected by the trigger only when QCD initial-state radiation (ISR) boosts the R-hadron pair system. Details of the event selection are given in Sec. VI.

VI. DATA SAMPLE AND EVENT SELECTION

The data sample used for this search was collected while tracking detectors, calorimeters, muon chambers, and magnets were operating normally and corresponds to an integrated luminosity of 3.2 fb$^{-1}$ in 2015. The luminosity measurement was calibrated during dedicated beam-separation scans, using the same methodology as that described in Ref. [55]. The uncertainty in the luminosity measurement is 5%.

Candidate events are selected by an $E_T^{\text{miss}}$ trigger, formed from energy deposits in the calorimeter [54], with a threshold of 70 GeV.

The event selection, described in detail below, was optimized relative to the $\sqrt{s} = 8$ TeV search [16] to improve the sensitivity to events with R-hadrons relative
to background events from SM processes. Major improvements include the addition of selections designed to reject tracks from hadrons and electrons, as well as changes to the isolation requirements, which now reject background events from overlapping tracks by identifying clusters in the Pixel or SCT detectors which are consistent with energy deposition from multiple particles. Two signal regions are defined: one targets metastable $R$-hadrons with shorter lifetimes, and the other is optimized for metastable or stable $R$-hadrons that are expected to pass through the muon spectrometer before decaying.

A primary vertex reconstructed from at least two well-reconstructed charged-particle tracks, each with $p_T > 400$ MeV, is required in order to remove noncollision background events. If there is more than one such vertex in an event, the primary vertex is defined as the vertex with the largest scalar sum of associated track momentum. To reduce backgrounds from SM processes, events must have $E_T^{\text{miss}} > 130$ GeV, where $E_T^{\text{miss}}$ is the off-line missing transverse momentum. Jets are reconstructed with the anti-$k_T$ algorithm [56] with radius parameter $R = 0.4$ using clusters of energy depositions in the calorimeter as inputs. Events are rejected if they contain a jet with $E_T > 20$ GeV that is consistent with noise contributions as determined based on timing and shower shape information. The off-line $E_T^{\text{miss}}$ is estimated from reconstructed electrons with $|\eta| < 2.47$, reconstructed muons with $|\eta| < 2.5$, calibrated jets with $p_T > 20$ GeV within $|\eta| < 4.5$, and reconstructed tracks which are not associated with other objects and that pass kinematic requirements designed to reject poorly reconstructed tracks.

For signal events, the trigger efficiency varies significantly depending on the model considered. For metastable $R$-hadrons with a mass of 1000 GeV, decaying into a 100 GeV neutralino and quarks, the trigger efficiency increases from 65% for $R$-hadrons with a lifetime of 50 ns to 95% for $R$-hadrons with a lifetime of 0.4 ns. For an $R$-hadron with a lifetime of 10 ns, the trigger efficiency increases slightly with the $R$-hadron mass, from 91% for a mass of 1000 GeV to 96% for a mass of 1600 GeV. For stable $R$-hadrons, the trigger efficiency is approximately 40% for all masses studied. All efficiencies quoted in this section include both the acceptance and the selection efficiency effects.

The off-line $E_T^{\text{miss}}$ calculation includes the momentum of identified muons, which is important to reject background events from SM processes with muons, such as $Z \rightarrow \mu \mu$, that pass the calorimetric $E_T^{\text{miss}}$ trigger. For events with metastable $R$-hadrons, the off-line $E_T^{\text{miss}}$ requirement is very efficient (around 95%) due to the neutralinos in the event. For stable $R$-hadrons that pass the trigger requirement, the off-line $E_T^{\text{miss}}$ cut is also efficient (around 90%), as the $R$-hadrons do not deposit significant energy in the calorimeter and in most events the $R$-hadron pair is boosted from initial-state radiation. Many $R$-hadrons fail to be identified as muons because they are in a neutral charge state going through either the inner detector or the muon spectrometer. Even if a charged $R$-hadron with $\beta < 1$ traverses the muon spectrometer, it can fail to be identified by the standard muon reconstruction algorithm as a prompt muon due to its late time of arrival. In a typical $R$-hadron model considered in this search, only 30% of stable $R$-hadrons with a reconstructed track in the inner detector are identified as prompt muons entering into the $E_T^{\text{miss}}$ calculation.

Selected events must contain at least one candidate $R$-hadron track. However, if more than one track in an event passes all requirements, only the highest-$p_T$ candidate is considered, in order not to bias the distribution of the variables and to allow for proper normalization of the background estimate.

For a track to be identified as a candidate, it must have $p_T > 50$ GeV and meet the following selection criteria. The track must contain at least seven clusters from silicon detector layers in order to make a good measurement of the track’s momentum, must have a cluster in the innermost Pixel layer if expected, and must be associated with the primary vertex. In order to ensure a good ionization measurement, the track must have at least two clusters in the Pixel detector used to measure $dE/dx$. These selection requirements on associated clusters are summarized as cluster requirements.

To reject background events with overlapping tracks that could produce clusters with significant measured ionization, the candidate track must not contain any cluster that is compatible with contributions from two or more tracks. This requirement is over 90% efficient for signal events with lifetimes of 3 ns or greater. For background tracks, it lessens the high ionization tail of the distribution from photon conversions and collimated particles, reducing the size of the background and leaving the remaining tracks with a $dE/dx$ distribution largely independent of background source. In order to reduce backgrounds from particles inside hadronic jets, the scalar sum of the $p_T$ ($\sum p_T$) of other tracks, with $p_T > 1$ GeV and consistent with the primary vertex, in a cone of size $\Delta R < 0.25$, where

$\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$, around the candidate track, must be less than 20 GeV.

The background from SM processes is further reduced by requiring that the candidate track momentum exceeds 150 GeV and that the relative uncertainty on the momentum measurement be less than 50%. The contribution from leptons from $W$ boson decays is reduced by imposing a condition that the transverse mass, $m_T$, built with the candidate track and the $E_T^{\text{miss}}$, be greater than 130 GeV, where

$$m_T = \sqrt{2p_T E_T^{\text{miss}}(1 - \cos(\Delta \phi(E_T^{\text{miss}}, \text{track}))} \quad (2)$$

and $\Delta \phi(E_T^{\text{miss}}, \text{track})$ is the azimuthal separation between the track and the $E_T^{\text{miss}}$. This selection has an efficiency between 78% and 89% for $R$-hadrons with lifetimes of 3 ns or greater, and reduces the background by roughly a factor of 3.
Background events due to tracks from electrons and hadrons are reduced by requiring that the energy of any jet associated with the candidate track must be less than the candidate track’s measured momentum. Jets are associated with tracks if they are within $\Delta R < 0.05$ of the candidate track and have $p_T > 20$ GeV. Background events containing electrons are further reduced by rejecting candidate tracks associated with a jet with 95% or more of its energy deposited in the electromagnetic calorimeter. These selections are found to be 95% efficient for all signal models studied.

Tracks identified as well-reconstructed muons with $p_T > 25$ GeV are rejected in the search for metastable $R$-hadrons with lifetimes of 30 ns or less. $R$-hadrons with a lifetime equal to or greater than 50 ns typically reach the muon spectrometer, and therefore no such veto is applied in the long-lived $R$-hadron search. To further reject hadronic background in the search for stable $R$-hadrons, the isolation requirement on the $\sum p_T$ of nearby tracks is tightened from 20 to 5 GeV for candidate $R$-hadron tracks, eliminating 30% of the remaining expected background.

The overall signal selection efficiency depends on the mass and lifetime of the signal, ranging from about 19% for a 1600 GeV $R$-hadron with a lifetime of 10 ns to less than 1% for $R$-hadrons with a lifetime of 0.4 ns. The trigger and off-line $E_T^{\text{miss}}$ requirements are more efficient for $R$-hadrons that decay before or inside the calorimeters. However, as the $R$-hadron lifetime decreases, the probability of reconstructing the track from the original $R$-hadron in the silicon detectors decreases. For these two reasons, $R$-hadrons with lifetimes in the 10–30 ns range have the highest selection efficiency.

The inefficiency for selecting $R$-hadrons with short lifetimes can be better understood by calculating the selection efficiency relative to a simple fiducial region that separates out acceptance effects due to basic kinematic properties of the $R$-hadrons. An event is considered to be within the fiducial acceptance if it contains at least one $R$-hadron that passes the following requirements at Monte Carlo generator level: charged at production, $p_T > 50$ GeV, $p > 150$ GeV, $|\eta| < 2.5$, and transverse decay distance $\geq 37$ cm from the origin. With this definition, the acceptance for $R$-hadron events with a lifetime of 1 ns ranges from 7% for a mass of 600 GeV to 4% for a mass of 1600 GeV. The selection efficiency for $R$-hadron events in the fiducial region has only a mild dependence on the mass; in the $R$-hadron mass range from 600 to 1600 GeV, the selection efficiencies are 20% to 25% ($T = 1$ ns), 35% to 39% ($T = 3$ ns), 39% to 45% ($T = 10$ ns), and 31% to 35% ($T = 30$ ns), respectively.

Table II summarizes the systematic uncertainties in the predicted signal yields. The systematic uncertainties are calculated and used independently for each mass bin and signal lifetime. In the table, only the maximum uncertainty is quoted.

### Table I

<table>
<thead>
<tr>
<th>Selection level</th>
<th>Expected signal events</th>
<th>Observed events in $3.2$ fb$^{-1}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Generated</td>
<td>26.0 $\pm$ 0.3</td>
<td></td>
</tr>
<tr>
<td>$E_T^{\text{miss}}$ trigger &amp; preselection</td>
<td>24.8 $\pm$ 0.3 (95%)</td>
<td></td>
</tr>
<tr>
<td>$E_T^{\text{miss}} &gt; 130$ GeV</td>
<td>23.9 $\pm$ 0.3 (92%)</td>
<td></td>
</tr>
<tr>
<td>Track $p_T &gt; 50$ and cluster requirements</td>
<td>10.7 $\pm$ 0.2 (41%)</td>
<td>368324</td>
</tr>
<tr>
<td>Isolation requirement</td>
<td>9.0 $\pm$ 0.2 (35%)</td>
<td>108079</td>
</tr>
<tr>
<td>Track $p &gt; 150$ GeV</td>
<td>6.6 $\pm$ 0.2 (25%)</td>
<td>47463</td>
</tr>
<tr>
<td>$m_T &gt; 130$ GeV</td>
<td>5.8 $\pm$ 0.2 (22%)</td>
<td>18746</td>
</tr>
<tr>
<td>Electron &amp; hadron veto</td>
<td>5.5 $\pm$ 0.2 (21%)</td>
<td>3612</td>
</tr>
<tr>
<td>Muon veto</td>
<td>5.5 $\pm$ 0.2 (21%)</td>
<td>1668</td>
</tr>
<tr>
<td>Ionization requirement</td>
<td>5.0 $\pm$ 0.1 (19%)</td>
<td>11</td>
</tr>
</tbody>
</table>
The trigger and off-line $E^\text{miss}_T$ selections are sensitive to the modeling of ISR in the event, in particular for signal events that do not decay inside the detector volume. Half of the difference between the selection efficiency of the PYTHIA 6 modeling of ISR in the event, in particular for signal events simulation. The uncertainty in the differences between efficiencies measured in data and in $Z\rightarrow\mu\mu$ samples and those reweighted with MG5_aMC@NLO is taken as an uncertainty due to radiative effects.

The $E^\text{miss}_T$ trigger modeling is studied in a sample of $Z\rightarrow\mu\mu$ events and an uncertainty is assigned based on differences between efficiencies measured in data and in simulation. The uncertainty in the $E^\text{miss}_T$ calibration comes primarily from uncertainties on the jet energy scale. A systematic uncertainty in the measurement of the Pixel ionization is evaluated by comparing the ionization of an inclusive sample of simulated and observed low-momentum tracks. An asymmetric uncertainty is assigned due to the selection efficiency difference observed after the most probable value of the signal distribution is reweighted based on the comparison of low-momentum tracks between data and simulation. The uncertainty in the signal acceptance is asymmetric since the ionization in simulation is systematically larger than that observed in data and larger at low mass (7.1% at 1000 GeV) than at high mass (2.0% at 1800 GeV). An uncertainty in the efficiency of the muon veto is estimated by studying the modeling of the muon spectrometer’s timing measurement: it is 3.2% for the 30 ns samples and negligible for the shorter lifetime samples, as the $R$-hadrons rarely reach the spectrometer.

The signal selection efficiency is found to be largely independent of the number of pileup events in simulation. However, to cover any remaining difference, efficiency differences from variations in the pileup distributions are accounted for as an additional source of uncertainty. The uncertainty on the momentum scale and resolution of the tracks results in a small uncertainty in the selection efficiency. Variations of the requirements to reject tracks from electrons and hadrons are found to have a negligible effect on the signal efficiency. The theoretical uncertainty in the signal cross section is described in Sec. IV.

### VIII. BACKGROUND ESTIMATION

The expected shape and normalization of the mass distribution of background events from SM processes is derived from data. The distributions of key variables are extracted in two control regions in data and these templates are used to generate the expected background distribution in the signal region. The choice of control samples takes into account the measured correlations between the key variables $p$, $dE/dx$ and $\eta$, and the control region selections minimize the possible signal contamination. The same regions are used in the searches for stable and metastable particles, except for the rejection of track candidates identified as muons in the latter case and for the tighter track isolation requirement in the former case.

The first control region (hereafter called CR1) is selected by applying all the selections described in Sec. VI, except for the requirement on high ionization, which is inverted. The tracks in this control region are kinematically similar to those in the signal region for SM processes, and the expected $p$ and $\eta(p)$ distributions of tracks from background processes are extracted from this region. The second sample (CR2 hereafter) is selected by inverting the $E^\text{miss}_T$ requirement ($E^\text{miss}_T < 130$ GeV), while keeping all other selections unchanged. As more than 90% of $R$-hadron events that pass the trigger have $E^\text{miss}_T > 130$ GeV, while the $E^\text{miss}_T$ distribution from SM background processes falls steeply in the range 70–130 GeV, CR2 is background dominated. Tracks from this region are used to derive $dE/dx$ templates in bins of $\eta$, as $E^\text{miss}_T$ and $dE/dx$ are uncorrelated for tracks from SM processes. The signal contamination is less than 1% in both control regions for all $R$-hadron masses and lifetimes considered in this search.

The shape of the background in the signal region is estimated from the control region templates with the following procedure: 1 million $(p, \eta, dE/dx)$ triplets are generated, where the momentum is randomly sampled according to the template from CR1 tracks, the pseudorapidity is generated according to the $\eta(p)$-binned templates based on CR1 tracks, and the ionization is sampled according to $dE/dx(\eta)$-binned templates from CR2 tracks. This procedure maintains the measured correlations between $p$, $\eta$, and $dE/dx$ for tracks from SM sources. The particle mass $M$ is calculated given the predicted $dE/dx$ and $p$ values, using the technique explained in Sec. III.

The normalization of the generated background is obtained by scaling the background to the data in the shoulder region of the mass distribution (i.e. $M < 160$ GeV), where a possible signal has already been excluded [16]. The normalization is performed on the samples before the ionization requirement.
and by changing the IBL ionization correction by regions, but requiring the track momentum to be in the same requirements as for the nominal signal and control validation regions. The validation regions are selected with Table III.

The complete procedure is tested on signal-depleted mass distributions of the background and the observed data, before the ionization requirement, are shown in Fig. 3 for the metastable R-hadrons selection.

Systematic uncertainties in the background estimate are evaluated by varying the relative fraction of muons in the control regions, using an analytical description to vary the shape of the high ionization tail in the dE/dx template, and by changing the IBL ionization correction by $\pm 1\sigma$. The statistical uncertainty in the template shape, which is estimated by Poisson fluctuations of the templates, dominates. There is also a small statistical uncertainty from the size of the normalization region. The breakdown of uncertainties in the background estimate is shown in Table III.

The background predictions show both the statistical and systematic uncertainty as well as the observed events for the metastable and stable R-hadron selection. The observed mass distribution in data is shown in Fig. 4 for both selection regions, along with the background expectation.

No evidence of a signal above the background is observed. Expected and observed upper limits on R-hadron production cross sections are evaluated at 95% CL for a discrete set of R-hadron mass values by counting simulated signal, background and observed data events that survive the selection in a $\pm 1.4\sigma$ wide mass window around the position of each simulated signal mass peak for the corresponding mass hypothesis being probed. The peak position and width are estimated by a Gaussian fit to each individual signal mass distribution ($\sigma \approx 280/[\approx 500]$ GeV at an R-hadron mass of 1200 [1600] GeV). When the lifetime of the R-hadron is long enough to reach the muon spectrometer, the stable selection becomes more efficient than the metastable one and improves the expected sensitivity. This transition happens around 50 ns. The choice of whether to apply the metastable or stable R-hadron selection to each lifetime is based on the best expected limit. Based on this procedure, the stable selection is applied to both the 50 ns and stable lifetimes.

The cross-section upper limits are extracted, separately for the stable and the metastable searches (in the latter case, for each lifetime value), with the $CL_S$ method [58], using the profile likelihood ratio as a test statistic. In the procedure, the systematic uncertainties in the signal and background yields, as evaluated in Sec. VII and Sec. VIII, are treated as Gaussian-distributed nuisance parameters.

The cross-section upper limits are extracted, separately for the stable and the metastable searches (in the latter case, for each lifetime value), with the $CL_S$ method [58], using the profile likelihood ratio as a test statistic. In the procedure, the systematic uncertainties in the signal and background yields, as evaluated in Sec. VII and Sec. VIII, are treated as Gaussian-distributed nuisance parameters.

### Table III

<table>
<thead>
<tr>
<th>Source of uncertainty</th>
<th>[%]</th>
</tr>
</thead>
<tbody>
<tr>
<td>Statistical uncertainty from control region templates</td>
<td>15</td>
</tr>
<tr>
<td>Statistical uncertainty from normalization region</td>
<td>3</td>
</tr>
<tr>
<td>Analytical description of dE/dx</td>
<td>4</td>
</tr>
<tr>
<td>Particle species composition of CR2</td>
<td>3</td>
</tr>
<tr>
<td>IBL ionization correction</td>
<td>4</td>
</tr>
</tbody>
</table>

### Table IV

<table>
<thead>
<tr>
<th>Selection region</th>
<th>Background expected</th>
<th>Data</th>
</tr>
</thead>
<tbody>
<tr>
<td>Metastable R-hadron</td>
<td>11.1 $\pm$ 1.7 $\pm$ 0.7</td>
<td>11</td>
</tr>
<tr>
<td>Stable R-hadron</td>
<td>17.2 $\pm$ 2.6 $\pm$ 1.2</td>
<td>16</td>
</tr>
</tbody>
</table>
The statistical uncertainty in the background distribution also takes into account the uncertainty due to the normalization. Lower limits on the $R$-hadron mass are then derived, for each lifetime, by comparing the measured cross-section upper limits to the theoretically predicted production cross section. The resulting lower limits set on the mass of the stable and metastable particles are summarized in Table V.

The strongest mass limits are obtained for lifetimes of 10 ns or more. The cross-section upper limits for 10 ns

<table>
<thead>
<tr>
<th>Mass [GeV]</th>
<th>0</th>
<th>500</th>
<th>1000</th>
<th>1500</th>
<th>2000</th>
<th>2500</th>
<th>3000</th>
</tr>
</thead>
<tbody>
<tr>
<td>Entries / 50 GeV</td>
<td>$3 \times 10^{-10}$</td>
<td>$2 \times 10^{-10}$</td>
<td>$1 \times 10^{-10}$</td>
<td>$1 \times 10^{-9}$</td>
<td>$1 \times 10^{-8}$</td>
<td>$1 \times 10^{-7}$</td>
<td>$1 \times 10^{-6}$</td>
</tr>
</tbody>
</table>

$\tilde{g} = 800 \text{ GeV}$, $\tau(\tilde{g}) = 3 \text{ ns}$
$\tilde{g} = 1600 \text{ GeV}$, $\tau(\tilde{g}) = 10 \text{ ns}$

$\tilde{g}$ is the gluino, a supersymmetric particle that can be produced in proton-proton collisions at the LHC. The mass limits are derived using the $R$-hadron mass, which is a measure of the mass of the gluino.

The observed 95% CL lower limit on mass is given in Table V. The metastable selection is more efficient for lifetimes of 50 ns or more.

<table>
<thead>
<tr>
<th>Selection</th>
<th>$\tau$ [ns]</th>
<th>$M_{\text{obs}}$ [GeV]</th>
<th>$M_{\text{exp}}$ [GeV]</th>
</tr>
</thead>
<tbody>
<tr>
<td>Metastable</td>
<td>0.4</td>
<td>740</td>
<td>730</td>
</tr>
<tr>
<td>Metastable</td>
<td>1.0</td>
<td>1110</td>
<td>1150</td>
</tr>
<tr>
<td>Metastable</td>
<td>3.0</td>
<td>1430</td>
<td>1470</td>
</tr>
<tr>
<td>Metastable</td>
<td>10</td>
<td>1570</td>
<td>1600</td>
</tr>
<tr>
<td>Metastable</td>
<td>30</td>
<td>1580</td>
<td>1620</td>
</tr>
<tr>
<td>Metastable</td>
<td>50</td>
<td>1540</td>
<td>1590</td>
</tr>
<tr>
<td>Stable</td>
<td>50</td>
<td>1590</td>
<td>1590</td>
</tr>
<tr>
<td>Stable</td>
<td>stable</td>
<td>1570</td>
<td>1580</td>
</tr>
</tbody>
</table>

The mass limits are derived using the $R$-hadron mass, which is a measure of the mass of the gluino. The metastable selection is more efficient for lifetimes of 50 ns or more.

$\tilde{g} \rightarrow q\bar{q} \tilde{\chi}_0$, $m(\tilde{\chi}_0) = 100 \text{ GeV}$

$\tilde{g}$ is the gluino, a supersymmetric particle that can be produced in proton-proton collisions at the LHC. The mass limits are derived using the $R$-hadron mass, which is a measure of the mass of the gluino.

The observed 95% CL lower limit on mass is given in Table V. The metastable selection is more efficient for lifetimes of 50 ns or more.

The strongest mass limits are obtained for lifetimes of 10 ns or more. The cross-section upper limits for 10 ns

$\tilde{g} = 13 \text{ TeV}$, 3.2 fb$^{-1}$

$\tilde{g}$ is the gluino, a supersymmetric particle that can be produced in proton-proton collisions at the LHC. The mass limits are derived using the $R$-hadron mass, which is a measure of the mass of the gluino.

The observed 95% CL lower limit on mass is given in Table V. The metastable selection is more efficient for lifetimes of 50 ns or more.

The strongest mass limits are obtained for lifetimes of 10 ns or more. The cross-section upper limits for 10 ns
lifetime $R$-hadrons decaying to $q\bar{q}$ plus a light neutralino of mass $m(\tilde{\chi}_0) = 100$ GeV are shown in Fig. 5. For this signal, the expected and observed lower mass limits are 1600 GeV and 1570 GeV, respectively. The 8 TeV result excluded $R$-hadrons with a 10 ns lifetime and masses up to 1185 GeV, using the lower edge of the $\pm 1\sigma$ band around the theoretical cross section.

The excluded regions on the lifetime-mass plane for $R$-hadrons decaying to $q\bar{q}$ plus a light neutralino of mass $m(\tilde{\chi}_0) = 100$ GeV are shown in Fig. 6.

**X. CONCLUSION**

This paper presents the results of a search for massive, long-lived particles with lifetimes of $O$(ns) or longer, produced in 3.2 fb$^{-1}$ of $pp$ collision data at $\sqrt{s}=13$ TeV at the LHC, and identified by anomalous specific ionization energy loss in the ATLAS Pixel detector. This analysis takes advantage of the addition of a fourth silicon layer to the Pixel detector, the increased parton luminosity, and analysis improvements to significantly extend a search performed at $\sqrt{s}=8$ TeV [16]. No excess of events is observed over the background estimate, and analysis improvements to significantly extend a search performed at $\sqrt{s}=8$ TeV [16].

**ACKNOWLEDGMENTS**

We thank CERN for the very successful operation of the LHC, as well as the support staff from our institutions without whom ATLAS could not be operated efficiently. We acknowledge the support of ANPCyT, Argentina; YerPhI, Armenia; ARC, Australia; BMWF and FWF, Austria; ANAS, Azerbaijan; SSTC, Belarus; CNPq and FAPESP, Brazil; NSERC, NRC and CFI, Canada; CERN; CONICYT, Chile; CAS, MOST and NSFC, China; COLCIENCIAS, Colombia; MSMT CR, MPO CR and VSC CR, Czech Republic; DNRF and DSNRC, Denmark; IN2P3-CNRS, CEA-DSM/IRFU, France; GNSF, Georgia; BMBF, HGF, and MPG, Germany; GSRT, Greece; RGC, Hong Kong SAR, China; ISF, I-CORE and Benoziyo Center, Israel; INFN, Italy; JSPS, Japan; CNRST, Morocco; FOM and NWO, Netherlands; RCN, Norway; MNiSW and NCN, Poland; FCT, Portugal; MNE/IFA, Romania; MES of Russia and NRC KI, Russian Federation; JINR; MESTD, Serbia; MSSR, Slovakia; ARRS and MIZŠ, Slovenia; DST/NRF, South Africa; MINECO, Spain; SRC and Wallenberg Foundation, Sweden; SERI, SNSF and Cantons of Bern and Geneva, Switzerland; MOST, Taiwan; TAEK, Turkey; STFC, U.K.; DOE and NSF, U.S. In addition, individual groups and members have received support from BCKDF, the Canada Council, CANARIE, CRC, Compute Canada, FQRNT, and the Ontario Innovation Trust, Canada; EPLANET, ERC, FP7, Horizon 2020 and Marie Skłodowska-Curie Actions, European Union; Investissements d’Avenir Labex and Idex, ANR, Région Auvergne and Fondation Partager le Savoir, France; DFG and AvH Foundation, Germany; Herakleitos, Thales and Aristeia programmes co-financed by EU-ESF and the Greek NSRF; BSF, GIF and Minerva, Israel; BRF, Norway; Generalitat de Catalunya, Generalitat Valenciana, Spain; the Royal Society and Leverhulme Trust, U.K. The crucial computing support from all WLCG partners is acknowledged gratefully, in particular from CERN and the ATLAS Tier-1 facilities at TRIUMF (Canada), NDGF (Denmark, Norway, Sweden), CC-IN2P3 (France), KIT/ GridKA (Germany), INFN-CNAF (Italy), NL-T1 (Netherlands), PIC (Spain), ASGC (Taiwan), RAL (UK) and BNL (U.S.) and in the Tier-2 facilities worldwide.

SEARCH FOR METASTABLE HEAVY CHARGED PARTICLES …

PHYSICAL REVIEW D 93, 112015 (2016)

SEARCH FOR METASTABLE HEAVY CHARGED PARTICLES …  PHYSICAL REVIEW D 93, 112015 (2016)

35c Department of Physics, Nanjing University, Jiangsu, China
35d School of Physics, Shandong University, Shandong, China
36 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Clermont-Ferrand, France
37 Nevis Laboratory, Columbia University, Irvington, New York, USA
38 Niels Bohr Institute, University of Copenhagen, København, Denmark
39a INFN Gruppo Collegato di Cosenza, Laboratori Nazionali di Frascati, Italy
39b Dipartimento di Fisica, Università della Calabria, Rende, Italy
40a AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
40b Marian Smoluchowski Institute of Physics, Jagiellonian University, Krakow, Poland
41 Institute of Nuclear Physics Polish Academy of Sciences, Krakow, Poland
42 Physics Department, Southern Methodist University, Dallas, Texas, USA
43 Physics Department, University of Texas at Dallas, Richardson, Texas, USA
44 DESY, Hamburg and Zeuthen, Germany
45 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
46 Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Dresden, Germany
47 Department of Physics, Duke University, Durham, North Carolina, USA
48 SUPA—School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
49 INFN Laboratori Nazionali di Frascati, Frascati, Italy
50 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
51 Section de Physique, Université de Genève, Geneva, Switzerland
52a INFN Sezione di Genova, Italy
52b Dipartimento di Fisica, Università di Genova, Genova, Italy
53a E. Andronikashvili Institute of Physics, Iv. Javakhishvili Tbilisi State University, Tbilisi, Georgia
53b High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
54 II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
55 SUPA—School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
56 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS/IN2P3, Grenoble, France
57 Department of Physics, Hampton University, Hampton, Virginia, USA
58 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, Massachusetts, USA
59 Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
60a Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
60b ZITI Institut für technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
61 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
62 Department of Physics, The Chinese University of Hong Kong, Shatin, N.T., Hong Kong, China
63 Department of Physics, The University of Hong Kong, Hong Kong, China
64 Department of Physics, The Hong Kong University of Science and Technology, Clear Water Bay, Kowloon, Hong Kong, China
65 Department of Physics, Indiana University, Bloomington, Indiana, USA
66 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
67 University of Iowa, Iowa City, Iowa, USA
68 Department of Physics and Astronomy, Iowa State University, Ames, Iowa, USA
69 Joint Institute for Nuclear Research, JINR Dubna, Russia
70 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
71 Graduate School of Science, Kobe University, Kobe, Japan
72 Department of Physics, Kyushu University, Fukuoka, Japan
73 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
74 Physics Department, Lancaster University, Lancaster, United Kingdom
75 INFN Sezione di Lecce, Italy
76 Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
77 Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
Faculty of Pure and Applied Sciences, and Center for Integrated Research in Fundamental Science and Engineering, University of Tsukuba, Tsukuba, Japan

Department of Physics and Astronomy, Tufts University, Medford, Massachusetts, USA

Department of Physics and Astronomy, University of California Irvine, Irvine, California, USA

INFN Gruppo Collegato di Udine, Sezione di Trieste, Udine, Italy

ICTP, Trieste, Italy

Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy

Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden

Department of Physics, University of Illinois, Urbana, Illinois, USA

Department of Physics, University of British Columbia, Vancouver, British Columbia, Canada

Department of Physics and Astronomy, University of Victoria, Victoria, British Columbia, Canada

Department of Physics, University of Warwick, Coventry, United Kingdom

Waseda University, Tokyo, Japan

Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel

Department of Physics, University of Wisconsin, Madison, Wisconsin, USA

Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany

Fakultät für Mathematik und Naturwissenschaften, Fachgruppe Physik, Bergische Universität Wuppertal, Wuppertal, Germany

Department of Physics, Yale University, New Haven, Connecticut, USA

Yerevan Physics Institute, Yerevan, Armenia

Centre de Calcul de l’Institut National de Physique Nucléaire et de Physique des Particules (IN2P3), Villeurbanne, France

Deceased.

Also at Department of Physics, King’s College London, London, United Kingdom.

Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

Also at Novosibirsk State University, Novosibirsk, Russia.

Also at TRIUMF, Vancouver BC, Canada.

Also at Department of Physics & Astronomy, University of Louisville, Louisville, KY, USA.

Also at Department of Physics, California State University, Fresno CA, USA.

Also at Department of Physics, University of Fribourg, Fribourg, Switzerland.

Also at Departamento de Física de la Universitat Autonoma de Barcelona, Barcelona, Spain.

Also at Departamento de Física e Astronomia, Faculdade de Ciencias, Universidade do Porto, Portugal.

Also at Tomsk State University, Tomsk, Russia.

Also at Universita di Napoli Parthenope, Napoli, Italy.

Also at Institute of Particle Physics (IPP), Canada.

Also at National Institute of Physics and Nuclear Engineering, Bucharest, Romania.

Also at Department of Physics, St. Petersburg State Polytechnical University, St. Petersburg, Russia.

Also at Department of Physics, The University of Michigan, Ann Arbor MI, United States of America.

Also at Centre for High Performance Computing, CSIR Campus, Rosebank, Cape Town, South Africa.

Also at Louisiana Tech University, Ruston LA, United States of America.

Also at Institutio Catalana de Recerca i Estudis Avancats, ICREA, Barcelona, Spain.

Also at Graduate School of Science, Osaka University, Osaka, Japan.

Also at Department of Physics, National Tsing Hua University, Taiwan.

Also at Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands.

Also at Department of Physics, The University of Texas at Austin, Austin TX, USA.

Also at Institute of Theoretical Physics, Ilia State University, Tbilisi, Georgia.

Also at CERN, Geneva, Switzerland.

Also at Georgian Technical University (GTU), Tbilisi, Georgia.

Also at Ochadai Academic Production, Ochanomizu University, Tokyo, Japan.

Also at Manhattan College, New York NY, USA.

Also at Hellenic Open University, Patras, Greece.

Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.

Also at School of Physics, Shandong University, Shandong, China.

Also at Moscow Institute of Physics and Technology State University, Dolgoprudny, Russia.

Also at Section de Physique, Université de Genève, Geneva, Switzerland.
hh Also at Eotvos Lorand University, Budapest, Hungary.
i Also at International School for Advanced Studies (SISSA), Trieste, Italy.
j Also at Department of Physics and Astronomy, University of South Carolina, Columbia SC, USA.
k Also at School of Physics and Engineering, Sun Yat-sen University, Guangzhou, China.
l Also at Institute for Nuclear Research and Nuclear Energy (INRNE) of the Bulgarian Academy of Sciences, Sofia, Bulgaria.
m Also at Faculty of Physics, M. V. Lomonosov Moscow State University, Moscow, Russia.
n Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.
o Also at National Research Nuclear University MEPhI, Moscow, Russia.
p Also at Department of Physics, Stanford University, Stanford CA, United States of America.
q Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.
r Also at Flensburg University of Applied Sciences, Flensburg, Germany.
s Also at University of Malaya, Department of Physics, Kuala Lumpur, Malaysia.
t Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.