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Nadolny, Jakub; Lara-Lopez, Maritza A.; Cervino, Miguel; Bongiovanni, Angel; Cepa, Jordi; de Diego, Jose A.; Perez Garcia, Ana Maria; Perez Martinez, Ricardo; Sanchez-Portal, Miguel; Alfaro, Emilio; Castaneda, Hector O.; Gallego, Jesus; Jesus Gonzalez, J.; Ignacio Gonzalez-Serrano, J.; Padilla Torres, Carmen P.; Pintos-Castro, Irene; Povic, Mirjana

Published in:
Astronomy & Astrophysics

DOI:
10.1051/0004-6361/201936205

Publication date:
2020

Document version
Publisher's PDF, also known as Version of record

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Citation for published version (APA):
The OTELO survey
Nature and mass-metallicity relation for Hα emitters at z ~ 0.4

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Received 28 June 2019 / Accepted 1 March 2020

ABSTRACT

Context. A sample of low-mass Hα emission line sources at z ~ 0.4 was studied in the context of the mass-metallicity relation (MZR) and its possible evolution. We drew our sample from the OSIRIS Tunable Emission Line Object (OTELO) survey, which exploits the red tunable filter of OSIRIS at the Gran Telescopio Canarias to perform a blind narrow-band spectral scan in a selected field of the Extended Groth Strip. We were able to directly measure emission line fluxes and equivalent widths from the analysis of OTELO pseudo-spectra.

Aims. This study aims to explore the MZR in the very low-mass regime. Our sample reaches stellar masses (M*) as low as 10^6 M☉, where 63% of the sample have M* < 10^7 M☉. We also explore the relation of the star formation rate (SFR) and specific SFR with M* and gas-phase oxygen abundances, as well as the M*-size relation and the morphological classification.

Methods. The M* were estimated using synthetic rest-frame colours. Using an χ2 minimization method, we separated the contribution of [N II]1658 to the Hα emission lines. Using the N2 index, we separated active galactic nuclei from star-forming galaxies (SFGs) and estimated the gas metallicity. We studied the morphology of the sampled galaxies qualitatively (visually) and quantitatively (automatically) using high-resolution data from the Hubble Space Telescope-ACS. The physical size of the galaxies was derived from the morphological analysis using GALAPAGOS2/GALFIT, where we fit a single-Sérsic 2D model to each source.

Results. We find no evidence for an MZR evolution from comparing our very low-mass sample with local SFGs from the Sloan Digital Sky Survey. Furthermore, the same is true for M*-size and M*-SFR relations, as we deduce from comparison with recent literature. Morphologically, our sample is mostly (63%) populated by late-type galaxies, with 13% of early-type sources. For the first time, we identify one possible candidate outlier in the MZR at z = 0.4. The stellar-mass, metallicity, colour, morphology, and SFR of this source suggest that it is compatible with a transitional dwarf galaxy.

Key words. galaxies: evolution – galaxies: formation – galaxies: fundamental parameters – galaxies: star formation – galaxies: structure

1. Introduction

The stellar mass (M*) and gas-phase metallicity (Z) are among the most fundamental properties of galaxies and tracers of galaxy formation and evolution. While the stellar mass provides information about the amount of gas that is locked up into stars, the metallicity gives indications on the history of star formation and the exchange of gas between an object and its environment.

The mass-metallicity relation (MZR) has been studied for several decades, beginning with the pioneering work of Lequeux et al. (1979). Since then, several authors have investigated the origin and behaviour of this relationship in low and
high redshifts, and in different environments and mass regimes. Tremonti et al. (2004) studied this relation for \( >5\times10^4 \) star-forming galaxies (SFGs) in a demonstration of the statistical power of the Sloan Digital Sky Survey (SDSS). They found a tight (\( \pm 0.1 \) dex) correlation between stellar mass and metallicity spanning over three orders of magnitude in stellar mass (down to \( 10^{6.5} M_\odot \)) and a factor of 10 in metallicity, with a median redshift of 0.1. Possible explanations of the origin of the MZR include (i) gas outflows. In this scenario, low-mass galaxies suffer a stronger effect because they have lower escape velocities and can more easily lose enriched gas through stellar winds than massive galaxies. This scenario was first proposed by Larson (1974; see also Spitoni et al. 2010; Tremonti et al. 2004). Furthermore, Belfiore et al. (2019) recently found that the inferred outflow loading factor decreases with stellar mass. (ii) The second possible explanation might be the downsizing effect: massive galaxies process their gas faster, on shorter timescales and at earlier epochs than low-mass galaxies, where the star formation is slower and extends over longer periods (e.g. Cowie et al. 1996; Thomas et al. 2010). (iii) Another possible mechanism for the MZR is the dependence of stellar mass or star formation rate (SFR) on the initial mass function (IMF; e.g. Köppen et al. 2007; Gunawardhana et al. 2011). However, it is still a matter of debate which of these mechanisms or which combination of them plays a more important role in the formation and evolution of galaxies.

With the advent of large surveys, there has been an effort to observe fainter and higher redshift objects because they would provide strong constraints on our understanding of how galaxies evolve. For instance, the Galaxy and Mass Assembly (GAMA; Driver et al. 2011; Liske et al. 2015) survey observed \( \sim300\,000 \) galaxies two orders of magnitude deeper than SDSS, covering an area of \( \sim 286 \) deg\(^2\). In the survey, an evolution of metallicity of \( \sim 0.1 \) dex was found for massive galaxies at redshift \( >0.35 \) (Lara-López et al. 2013). At higher redshifts, the VIMOS Very Large Telescope (VLT) Deep Survey (VVDS; Garilli et al. 2008) provides spectroscopic data of \( >30 \) 000 galaxies up to redshift \( z \approx 6.7 \) over a total of \( >8 \) deg\(^2\) down to \( i_{AB} > 24.75 \). Using VVDS data, Lamarre et al. (2009) studied the MZR up to \( z \approx 0.9 \). They defined two samples: a wide (\( i_{AB} < 22.5 \)) and a deep sample (\( i_{AB} < 24 \)). They used three different metallicity calibrators depending on the redshift range. All methods were then normalized to the Charlot & Longhetti (2001) calibrator, which is in agreement with the Tremonti et al. (2004) metallicities. Assuming that the MZR shape is constant with redshift, they found a stronger metallicity evolution in the wide sample and concluded that the MZR is flatter at higher redshift. They arrived at the same conclusion for the luminosity-metallicity relation (LZR) in the wide sample.

There is an incentive to also analyse the low-mass end of the MZR because flattening in this relation has been reported. For instance, jimmy et al. (2015) found a flattening of the MZR relation using a sample of local (\( D < 20 \) Mpc) dwarf galaxies with the VIMOS integral field unit (IFU) spectrograph on the VLT. They also found a clear dependence of the MZR and LZR on SFR and H\( \text{I} \)-gas mass. They used this finding to explain the observed scatter at the low-mass (\( < 10^8 M_\odot \)) end of these relations. Using a sample of 25 dwarf galaxies, Lee et al. (2006) also extended the MZR for galaxies down to a \( M_\star \sim 10^8 M_\odot \), roughly 2.5 dex lower in stellar masses than Tremonti et al. (2004). Using Spitzer mid-IR (MIR) (4.5\( \mu \)m) photometry, they found a comparable scatter in the MZR over 5.5 dex in stellar mass. They concluded that if galactic winds are responsible for the MZR, increasing dispersion is expected, as they observed in their analysis.

A clear extreme in the MZR is given by extremely metal-deficient (XMD) galaxies (\( \leq 0.1 Z_\odot \)). As found by Thuan et al. (2016), XMD galaxies show a clear trend of increasing gas mass fraction with decreasing metallicity, mass, and luminosity. On the same line, Sánchez Almeida et al. (2014, 2015) have argued that the low metallicities of the XMD galaxies are an indicator of infall of pristine gas. However, Thuan & Iztov (2005) found that low metallicity alone cannot explain the hard ionising radiation observed in blue compact dwarf (BCD) galaxies, and proposed fast radiative shocks (velocity \( > 450 \) km s\(^{-1} \)) associated with supernovae explosions of massive Population III stars. Ekta & Chengalur (2010) examined the MZR of XMD galaxies and compared it with the MZR of BCD and dwarf irregular galaxies. They proposed that XMD galaxies are extremely metal depleted due to better mixing of the inter-stellar medium (ISM).

The fact that there are outliers of the MZR makes the situation even more interesting. Peebles et al. (2008) found a sample of \( 41 \) local (\( z < 0.048 \)), low-mass, high-oxygen abundance outliers from the Tremonti et al. (2004) SDSS-based relation. They argued that these compact, isolated, and morphological undisturbed sources may be transitional dwarf galaxies. Their red colours (approaching or on the red sequence) and low gas fraction suggest that they are in the final stage of their star formation period. Moreover, using a larger sample of galaxies from the SDSS and DEEP2 surveys with stellar masses down to \( 10^7 M_\odot \), Zahid et al. (2012) found that the properties of low-mass metal-rich galaxies (the outliers from MZR) are, in agreement with previous works, possible transitional objects between gas-rich dwarf irregular and gas-poor dwarf spheroidal and elliptical galaxies.

In consideration of previous studies, we wish to shed light on the low-mass end of the MZR. To this aim, we used data from the OSIRIS Tunable filter Emission-Line Object (OTELO) survey reported by Bongiovanni et al. (2019, hereafter OTELO-I) to study a sample of low-mass H\( \alpha \) emission line sources (ELS) at redshift \( \sim 0.4 \), firstly described in Ramón-Pérez et al. (2019, hereafter OTELO-II). As demonstrated in Bongiovanni et al. (2020), OTELO has allowed validation of the narrow-band (NB) scan technique using tunable filters (TFs) for finding faint populations of SFGs, which makes this survey the most sensitive to date in terms of minimum emission-line flux (\( \sim 5 \times 10^{-19} \) erg s\(^{-1} \) cm\(^{-2} \)) and observed equivalent width (EW \( \approx 6 \) Å).

This article is organized as follows. In Sect. 2 we present the data and sample selection process together with its characteristics and the external catalogues we used. In Sect. 3 we describe the methods for estimating stellar mass, gas metallicity, SFR, and morphology determination from OTELO data for the science case of this study. In Sect. 4 we present our analysis and discussion, followed with our conclusions in Sect. 5. When necessary, we adopt the cosmology from Planck Collaboration XIII (2016) with \( \Omega_\Lambda = 0.69, \Omega_m = 0.31 \), and \( H_0 = 67.8 \text{ km s}^{-1} \text{ Mpc}^{-1} \).

2. Data and sample selection

We study a sample of H\( \alpha \) ELSs with stellar masses in the range of \( 10^8 < M_\star / M_\odot < 10^{10} \), which places our sample in the low-mass regime. To this aim, we used data products of OTELO. A brief description follows.

The core of the OTELO survey consists of red tunable filter (RTF) observations, which were carried out in a total of 108 dark hours with the OSIRIS (Cepa et al. 2000) instrument at the 10.4 m Gran Telescopio Canarias\(^1\) (GTC). These observations

\(^1\) http://www.gtc.iac.es/gtc/gtc.php
cover a \(-56\) arcmin\(^2\) section of the Extended Groth Strip (EGS) region in six dithered pointings. The nominal central wavelengths at the centre of the field of view (FOV) range from 9070 to 9280 Å with \(\Delta\lambda\) of 6 Å steps. A full width at half-maximum (FWHM) of 12 Å was adopted for each one of the resulting 36 RTF tunings. This sampling of the RTF is optimal for deblending the H\(\alpha\) and [N\(\text{II}\)]\(6583\) lines (Lara-López et al. 2010a) for systems at \(z \sim 0.4\). In addition, all RTF images were combined in a single OTELO\(_\text{deep}\) image. A source detection on this image was performed using the SExtractor software (Bertin & Arnouts 1996), providing 11 237 sources. The integrated flux (OTELo\(_\text{deep}\)) of the sources detected in this image has a limiting magnitude of 26.4 [AB] (50% completeness).

The OTELO\(_\text{deep}\) image was also used to measure the point spread function (psf)-model fluxes of these 11 237 sources on registered and resampled public images in the optical and near-infrared (i.e. \(u, g, r, i, z, J, H\) and \(K\) bands) obtained from the AEGIS collaboration\(^2\). These data were cross matched in X-rays, far-UV (FUV) and near-UV (NUV)\(^3\), as well as those from Herschel and Spitzer IR bands to obtain additional photometric data for the sources so correlated. In addition, high-resolution \(F606W\) and \(F814W\) image stamps from the ACS of the Hubble Space Telescope (HST) of each source, when found, were obtained for further studies, in particular for morphological classification.

Related to the scope of this work, the resulting spectral energy distributions (SED) for each OTELO source were processed with the LePhare\(^4\) code (Arnouts et al. 1999; Ilbert et al. 2006) using a galaxy template library with the four standard Hubble types E/S0, Sbc, Scd, and Irr from Coleman et al. (1980) and six SFGs templates from Kinney et al. (1996), and an active galactic nucleus/quasi-stellar object (AGN/QSOs) library with two Seyfert, three type-1 and two type-2 QSOs and three composite (starburst+QSO) templates created by Polletta et al. (2007). This procedure allowed us to obtain photometric redshift solutions \(z_{\text{phot}}\), which were independently calculated including and excluding the OTELo\(_\text{deep}\) photometry in the input SED.

This provided a set of possible \(z_{\text{phot}}\) solutions, along with the two best-fit SED templates and the corresponding extinction \(E(B-V)\) estimates. The final best solution was obtained by the analysis of the best \(\chi^2\) solution obtained together with the estimated redshift error, \(\sigma_{z_{\text{phot}}}\), defined by

\[
\sigma_{z_{\text{phot}}} = |Z_{\text{BEST_LOW_deepX}} - Z_{\text{BEST_HIGH_deepX}}|, \tag{1}
\]

where \(Z_{\text{BEST_LOW_deepX}}\) and \(Z_{\text{BEST_HIGH_deepX}}\) are the 1\(\sigma\) deviations of the estimated redshift for the given object, while \(X\) stands for \(Y\) or \(N\), including or excluding the OTELO\(_\text{deep}\) photometry in the SED fitting, respectively. Finally, the astrometric and photometric data for each source, along with the photometric redshift estimates and the redshift data from other surveys DEEP2 (Newman et al. 2013) and CFHTLS T0004 Deep3 (Coupon et al. 2009), were included in the dubbed OTELo multi-wavelength catalogue.

The RTF scan data were also used to provide a pseudo-spectrum (PS) for each entry in the OTELO multi-wavelength catalogue. The PS are properly calibrated in flux and wavelength and differ from conventional spectra in that the former are affected by the distinctive RTF transmission profile. Such PS typically contain 36 wavelength points (\(R \sim 700\)), but the actual number of points and the wavelength range covered by each PS depend on the position of the object in the FOV. Some PS examples are given in Fig. 1. For each PS a pseudo-continuum flux \(f_c\), which itself is defined as the median flux of the PS, and a deviation around \(f_c\) \(\sigma_{f_{\text{mod},c}}\) were estimated. An object was qualified as an ELS candidate when either two adjacent points of the PS were above \(f_c + 2\sigma_{f_{\text{mod},c}}\) or only one was above it but it has an adjacent point above \(f_c + \sigma_{f_{\text{mod},c}}\). With these criteria, a list 5322 preliminary ELS candidates for the whole survey was obtained. Further details about the survey design, observations, data reduction, photometric redshift estimations, the construction of the multi-wavelength catalogue, the PS preparation, and the preliminary ELS selection are given in OTELO-I.

2.1. H\(\alpha\) + [N\(\text{II}\)] sample selection

We used the OTELO data described above to select the H\(\alpha\) sample as follows:

(a) According to OTELO-II, we selected the sources where the best \(z_{\text{phot}}\) (including and excluding the OTELO\(_\text{deep}\) photometry in the SED)
photometry) was in the range $0.2 < z_{\text{phot}} < 0.5$ with the further constraint that $\sigma_{\text{phot}} < 0.2 \times (1 + z_{\text{phot}})$. The total number of sources in this redshift window is 2289. In this selection we lost some Hα + [N II] candidates whose best $z_{\text{phot}}$ would be obtained with AGN/QSOs libraries instead of galaxy templates. This does not affect the goals of this work. In addition, we are also aware that our selection includes possible [S II] $\lambda\lambda6716,6731$ at $z \sim 0.3$. The best-fit templates correspond to Scd, Im, or SFGs types, with the exception of two sources where the best $z_{\text{phot}}$ was provided by an E/S0 galaxy type. Within this set of 2289 sources, we found 73 preliminary ELS candidates to Hα + [N II].

(b) We visually inspected the 73 candidates using our web-based graphic user interface (GUI)\(^5\) where all data about a user-selected source can be displayed; in particular, it includes the PS, the thumbnails of the images in available bands, and the SExtractor segmentation maps together with high-resolution HST images, $z_{\text{phot}}$ solution including and excluding OTELO\(_{\text{deep}}\) photometry obtained from our photo-z analysis and their comparison with the input photometry, the $z_{\text{phot}}$ solution from the CFHTLS T0004 Deep3 photo-z catalogue, and the $z_{\text{spec}}$ from DEEP2.

The high-resolution HST-ACS images allow us to identify possible contamination of the PS from neighbouring sources that are not resolved in the ground-based images. The comparison of the best-fit SED and the source data, together with the additional redshift estimates, provides us a confidence in our $z_{\text{phot}}$ solution. Finally, the PS allows us to distinguish between Hα + [N II] SFG sources, [S II]$\lambda\lambda6716,6731$, or AGN galaxies with Hα + [N II] emission. We took advantage of the fact that the [N II] $\lambda6583$ over Hα ratio must be lower than 0.4 (using the extreme case used by Pettini & Pagel 2004, see below) in SFGs. In addition, the [S II]$\lambda\lambda6716$ over [S II]$\lambda6731$ line ratio must range between 1.49 and 0.44 (see McCall 1984; Cid Fernandes et al. 2013 in the context of tunable filters). This means that any doublet where the redward line ([N II]$\lambda6583$ or [S II]$\lambda6731$) is stronger than about half of the blueward line (Hα or [S II]$\lambda6716$) should be a Hα + [N II] AGN or a [S II]$\lambda\lambda6716,6731$ source. This process reduced the sample of 73 candidates to a sample with 18 sources where Hα and [N II]$\lambda6583$ can be measured. In addition, we also obtained a $z_{\text{GUESS}}$ value (from the web-based GUI), which was obtained by the association of the brightest RTF slice and the Hα emission line at a given observed wavelength.

(c) As a final cross-check, we searched for objects with DEEP2 spectroscopic redshift in the range of $0.35 \leq z_{\text{spec}} \leq 0.42$, which covers the range where Hα line would be detected in the OTELO spectral window. This redshift range would overlap with [S II]$\lambda\lambda6716,6731$ sources (0.32 $\leq z \leq 0.38$). DEEP2 data provides 18 sources in common with our selection in point (b), 9 of them do not show any emission feature in the OTELO PS, and 7 of them are in common with our selection from point (b). The remaining 2 objects had a best $z_{\text{phot}}$ solution for the galaxy library outside of our redshift selection range (although one of them provides a $z_{\text{phot}}$ solution in the Hα + [N II] range using the AGN/QSOs library with a Sy-2 galaxy template). After visual inspection of the PS, we verified that one source of this set is indeed a Hα + [N II] emitter (id: 625) and the second one displays only the [N II]$\lambda6583$ line (and [S II]$\lambda\lambda6716,6731$ lines on the red side, id: 6059) in its PS. We therefore included only id: 625 source in the final sample, which increases to 19 objects. We note that this additional source was not included in the preliminary ELS list because the emission line is truncated in one extremity of the PS (see the middle panel of Fig. 1), but it was selected as a Hα + [N II] SFG candidate with the colour-magnitude ELS selection technique that is traditionally used in deep NB surveys (Pascual et al. 2007). In our case, a (z-OTELO$_{\text{int}}$) colour was obtained for all the OTELO sources detected in z band. See OTELO-I for details about this ELS selection procedure.

In summary, the final sample of Hα ELSs SFG candidates consists of 19 sources; we refer to this hereafter as the Hα sample. Further refinements (i.e. the flux and metallicity estimations) are described in the following sections. The examples given in Fig. 1 show a PS for a Hα + [N II] source (id: 2146), a source with a truncated Hα (id: 625), and [S II]$\lambda\lambda6716$ emission and a [S II]$\lambda\lambda6716,6731$ source (id: 3289, not in the Hα sample). The top part of Table 1 shows a summary of the number of sources at different steps of the selection process. In the appendix we show figures with the PS of the sources that compose the Hα sample.

### 2.2. Data from previous surveys

In order to compare our results with those from previous studies, we used the SDSS-Data Release 7 (SDSS-DR7, Abazajian et al. 2009) and the GAMA survey (Driver et al. 2011). From both datasets we considered only the SFGs selected through the BPT diagram (Baldwin et al. 1981) using the criteria of Kauffmann et al. (2003). In both surveys, the stellar masses were estimated using SED templates from Bruzual & Charlot (2003) and a Chabrier (2003) IMF. SDSS-DR7 fit templates to the observed SEDs, while stellar masses in GAMA were estimated using $(g-i)$ colours and a mass–luminosity relation (Taylor et al. 2011).

We selected the following sub-samples from the SDSS-DR7 (local and high-z) and the GAMA surveys: (i) A local SDSS sample with 0.005 $< z < 0.1$ and a signal-to-noise ratio ($S/N$) $> 3$ for Hα and $S/N > 4$ for the [N II] emission line. (ii) A higher

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redshift sample (high-z SDSS and GAMA) with 0.3 < z < 0.33 (upper-z limit due to the Hα line visibility in both surveys). The emission line S/N imposed for both high-z samples was the same as in (i).

Furthermore, to ensure completeness for the higher redshift samples, we imposed a cut in stellar mass of 10^{10} M_\odot. This cut is given by the limit in completeness of SDSS (see Weigel et al. 2016), and it approximately corresponds to an i-band absolute magnitude limit of M_\text{i} \sim -21 [AB]. These selection processes guarantee that we used robust and complete samples from both surveys.

We also used the data from publicly available catalogues of the VVDS\footnote{http://cesam.lam.fr/vvds/index.php}, in particular, from VVDS-Deep and VVDS-Ultra-Deep (VVDS-F0226-04). Data used from both surveys were selected as follows: redshift quality flag 1 < ZFLAG < 11 (this removes possible AGN from the sample, as described in Le Fèvre et al. 2013); the redshift range of 0.3 (Deep)/0.2 (Ultra-Deep) < z < 0.42, we decided to use z = 0.2 as the lower redshift limit for Ultra-Deep in order to increase the final number of sources.

3. Methods

3.1. Stellar masses

Stellar masses M_\star for the overall OTELO sample were computed using the López-Sanjuan et al. (2019; hereafter LS18) mass-to-light ratio for quiescent and SFGs (their Eqs. (10) and (11), respectively) employing rest-frame g- and i-band magnitudes, together with absolute i-band magnitude M_\text{i}. The rest-frame photometry was computed by shifting the SED_{\text{rest}} (see Sect. 2) to z = 0, and calculating the flux below the transmission curves for g, r, i, z (CFHTLS) and for the HST-ACS F606W and F814W filters. To estimate M_\star we used our z_{\text{OTELO}} (see Sect. 3.2). A second-order correction accounted for the Galactic dust extinction towards the EGS using the Galactic dust map from Schlafly & Finkbeiner (2011). The obtained stellar masses M_\star for the H\alpha sample are in the range of \sim 10^{6.8} < M_\star < 10^{10} M_\odot. This places our sample in a very low-mass regime where 63% of our sample has M_\star < 10^9 M_\odot. To test its consistency, we also calculated M_\star using the Taylor et al. (2011; hereafter T11) mass-to-light ratio. Both methods agree, with the most noticeable difference for blue SFGs, for which LS18 provides an additional quadratic factor that depends on the (g-i) colour. In Fig. 2 we compare the computed stellar masses using prescriptions from LS18 and T11 for all the OTELO sources (median value of M_\star = 10^{9.8} M_\odot) and for the H\alpha sample (median value of M_\star = 10^{8.4} M_\odot). In Fig. 3 we show the obtained stellar masses as a function of rest-frame (g-i) colour for the overall OTELO and the H\alpha ELS samples. For comparison, in the same figure, we show the SDSS overall sample together with the blue and red cloud empirical colour division of Bluck et al. (2014):

\begin{equation}
(g - r) = 0.06 \times \log M_\star - 0.01,
\end{equation}

obtained using data from SDSS-DR7. After we identified blue and red clouds in SDSS using the (g-r) colour, we represented them in terms of (g-i). Because it is only for illustrative purposes, we did not correct the blue and red cloud selection in terms of colour used.

Furthermore, we compared the obtained masses with the COSMOS2015 catalogue (Laigle et al. 2016) finding good agreement: our mass and colour estimates fall within the COSMOS2015 imprint. For the sake of clarity, we decided to omit this in Fig. 3, where a similar comparison is shown using data from SDSS-DR7. The resulting stellar masses of the sample we used and their associated uncertainties are discussed in Appendix B and shown in the fourth column of Table B.1.
3.2. Deconvolution of PS, metallicity, SFR, and sSFR

To obtain gas-phase metallicities and AGNs contamination, we used the N2 index:

\[
N2 = \log \frac{f([\text{NII}]+6583)}{f(H\alpha)},
\]

which requires a flux estimate. To this end, we used the so-called inverse deconvolution of the PS. In short, we assumed rest-frame model spectra defined by a constant continuum level, and Gaussian profiles of the [N II]+6548, H\alpha, and [N II]+6583 lines defined by their amplitude and a common line width for the three lines. The model to be compared with observational data is thus defined by the parameters \(f_{\text{mod}}(\lambda, f_{\alpha}, \sigma_{\text{line}}, f_{\text{NII}}, f_{\text{Halpha}}, f_{\text{NII}}+6583)\). We performed 10^6 Monte Carlo simulations for each object and varied the different parameters. The different parameters varied in following ranges. The redshift was varied in the range \(z_{\text{GUESS}} \pm 0.001\). The line width \(\sigma_{\text{line}}\) varied from 20 to 500 km s^{-1} rest-frame values, which translates into a range of 0.43–11 Å in spectral units at the H\alpha wavelength. The continuum \(f_{\alpha}\) level varied in the range \(f_{\text{mod}} \pm \sigma_{\text{mod}}\) with \(f_{\text{mod}}\) the median of the PS and \(\sigma_{\text{mod}}\) the root mean square around \(f_{\text{mod}},\) as defined in Sect. 2.1. The rest-frame \(f_{\text{Halpha}}\) amplitude varied between \(f_{\text{mod}}\) and the maximum value observed in the PS, \(f_{\text{PS}}\), taking into account the observational error \(\sigma(f_{\text{PS}})\), and taking into account the assumed Gaussian profile. This is given by Eq. (A.5). The rest-frame \(f(\text{NII}+6583)\) varied between 0.006 \(\cdot f_{\alpha}\) and the maximum possible value given by Eq. (A.5). Finally, the rest-frame \(f(\text{NII}+6583)\) was fixed to 1/3 \(\cdot f_{\text{NII}+6583}\) (Osterbrock & Ferland 2006). See Appendix A for details.

Each model was convolved with the OTELO filter response to produce a synthetic PS that was compared with the observed PS, and the corresponding \(\chi^2\) was obtained. The Monte Carlo simulations are designed to sample the multiparametric probability density function (PDF) from which modal values (i.e. best-fit solutions) and different confidence intervals around the modal value were obtained for each parameter (see details in Appendix A). In the following we consider error bars as the confidence intervals, which includes 25% and 68% (which would be the equivalent of \(\pm 1\sigma\) for the Gaussian case) of the total area of the PDF around the modal value as an estimate of the errors. The redshift value obtained for each source after deconvolution (hereafter \(z_{\text{OTELO}}\)) are given in Table B.1. Formally, the common maximum uncertainty of these values is about 0.001(1 + \(z_{\text{OTELO}}\)), as established in Bongiovanni et al. (2020).

We note that the possible flux-related parameters we obtained are strongly correlated, that is, each model has its own N2 index (i.e. metallicity), line equivalent widths, and so on, because the variation in one of the parameters of the model non-trivially affect the possible best solutions of the other parameters. In this way, we also obtained the corresponding PDFs of all the associated parameters (equivalent widths, N2 indices, and oxygen abundances). For all these inferred quantities we performed posteriori operations over the entire Monte Carlo set associated with each particular object. The corresponding PDFs (and associated errors) of the particular quantities can be obtained as shown in Appendix A. In the case of \(f_{\alpha}\), we corrected the model values for possible stellar absorption of underlying older components following Hopkins et al. (2003, 2013) using a EW_{\text{c}} correction value of 2.5 Å:

\[
F_{\text{cor}}^{H\alpha} = \frac{EW_{\text{Halpha}} + EW_{\text{c}}}{EW_{\text{Halpha}}} \times F_{\text{mod}}^{H\alpha},
\]

where \(F_{\text{mod}}^{H\alpha}\) and \(F_{\text{cor}}^{H\alpha}\) are the model before and after stellar absorption correction, respectively, and EW_{\text{Halpha}} refers to the equivalent width of the model before correction.

Finally, we computed the gas-phase oxygen abundance from the N2, with the absorption corrected \(F_{\text{cor}}^{H\alpha}\) value by means of Eq. (1) of Pettini & Pagel (2004) for all simulations:

\[
12 + \log (O/H) = 8.90 + 0.57 \times N2.
\]

Nominal values and errors for metallicities were calculated by analysing the metallicity PDF obtained from the Monte Carlo simulations, but we did not take the intrinsic dispersion in the calibration into account, which would be ±0.41 ±0.18 dex at 95% (68%) confidence interval (Pettini & Pagel 2004). Table B.1 shows the resulting \(f_{\text{Halpha}},\) \(EW_{\text{Halpha}}\) (both including the correction for stellar absorption), \(f(\text{NII}+6583),\) and \(12 + \log(O/H)\) (without calibration uncertainty) values with their 68% confidence intervals inCols. 6, 8, 7, and 5, respectively.

We note that by construction, the inverse deconvolution process always provides a solution for the intensity of the \(f(\text{NII}+6583)\) value, which would correspond to a 12 + log(O/H) = 7.63 (before the stellar absorption correction). These values would be below the OTELO line flux limit of 5 \times 10^{-19} \text{erg}^{-1} \text{cm}^{-2}\) as occurs for five objects (see Bongiovanni et al. 2020, for the flux limit estimation). These cases translate into a large uncertainty in the resulting N2 and metallicity values. We refer to these objects as low-flux sources and can only provide a \(f(\text{NII}+6583)\) (hence metallicity) limit.

In addition, the Pettini & Pagel (2004) method is valid in the range of \(-2.5 < N2 < -0.3\), which corresponds to the 7.47 to 8.73 12 + log(O/H) range. Values of 12 + log(O/H) equal to 8.67 would correspond to \(N2 = -0.4\), which defines an empirical division between SFGs and AGNs (Stanislawski et al. 2006). In order to remove possible AGNs from our sample, we used the same method as in Ramón-Pérez et al. (2019, as for narrow-line AGN), that is, the Cid Fernandes et al. (2010) diagnostic diagram (\(EW_{\text{Halpha}}\) vs. N2) shown in Fig. 4. In Fig. 4, the green and blue dashed lines are the limits of the empirical metallicity estimations from Yin et al. (2007) and Pettini & Pagel (2004) of −0.5 and −0.3, respectively. In this way, we found three sources as possible AGN and excluded them from the further analysis. The lower part of Table 1 shows the summary of sources of different classes after the inverse deconvolution and AGN discrimination process.

Finally, we corrected the H\alpha fluxes for dust attenuation using the reddening values obtained from LePhare SED fitting (quoted in the second column of Table B.1), and the empirical relation by Calzetti et al. (2000) estimated for galaxies at \(z = 0.5\) by Ly et al. (2012). We note that this is a lower value of \(L(\text{H}\alpha)\) because the extinction is obtained from the fit of galaxy templates whose intrinsic extinction is unknown. The derived values of \(E(B-V)\) are in the range of 0.0 and 0.3 mag, which is consistent with the extinction expected for low-mass galaxies, as shown in previous studies (e.g. Domínguez Sánchez et al. 2014). We also note that LePhare does not provide an evaluation of the uncertainty in \(E(B-V)\), therefore the uncertainty in this term was not taken into account to evaluate the uncertainty in \(L(\text{H}\alpha)\) (hence in the SFR, see Appendix B) The H\alpha line fluxes were used to derive SFRs following the standard calibration of Kennicutt (1998), but including the correction for a Chabrier (2003) IMF (see Kennicutt & Evans 2012, their Table 1), which is the one used in this work:

\[
\text{SFR}(M_{\odot} \text{yr}^{-1}) = 5.37 \times 10^{-42} L(\text{H}\alpha)[\text{erg s}^{-1}].
\]
The specific SFR (sSFR) was subsequently estimated as $SFR/M_*$. The resulting SFR and its associated uncertainty are listed in Col. 9 in Table B.1.

### 3.3. Morphology

Because the galaxies we analysed are some orders of magnitudes fainter than those in other surveys, it is mandatory to characterize this population as well as possible. Therefore, we performed a (i) quantitative and (ii) qualitative morphological classification. In both cases we decided to base this on the publicly available\(^7\) high-resolution HST-ACS ($F606W$ and $F814W$) data.

(i) We used the GALAPAGOS-2 software (Häussler et al. 2007) that was developed within the MegaMorph project\(^8\). GALAPAGOS-2 is an automatic way to provide the morphological parameters with SExtractor (Bertin & Arnouts 1996) and GALFIT (Peng et al. 2002). In this work we provide a single-Sérsic (1968) fit to sources detected in HST-ACS-$F814W$ using SExtractor in high-dynamical range (HDR) dual mode. The HDR mode maximizes the detection of faint sources, and the dual mode ensures that we measured the same sources in all the filters. As indicated in Peng et al. (2010), the formal uncertainties from GALFIT are only lower estimates. After simulations and tests of this software, Häussler et al. (2007) provided uncertainty estimates based on the comparison of input and output values as a function of surface brightness. Because we do not provide simulations like this, we cite the formal uncertainties from GALFIT, but we caution that these are the lower estimates. In Fig. 5 we show the uncertainties on effective radius and Sérsic index as a function of surface brightness defined following Häussler et al. (2007). Even if these are the lower limit errors, we can see the expected tendency of errors to increase with surface brightness. The resulting Sérsic index and its uncertainty is shown in Col. 10 of Table B.1.

(ii) Visual classification was made using MorphGUI, which is a graphic user interface for morphological classification developed by CANDELS (see Kartaltepe et al. 2015). We modified the interface in order to provide additional morphological classes to extend the classical Hubble scheme, to be able to take into account peculiarities found at higher redshifts. Our classification scheme is based on the work of Elmegreen et al. (2007): point-like (PL), spheroidal (Sph), disc (D), spheroidal-disc (Sph+D), tadpole (T), chain (C), clumpy cluster (CC), doubles (DD), and unclassified (Unclass).

The OTELO GTC/OSIRIS data have considerably lower resolution than HST-ACS (0.256 and 0.03 arcsec px\(^{-1}\), respectively), therefore we performed the match of the two catalogues by allowing multiple matches when there was more than one object in the HST-ACS $F814W$ image inside the Kron ellipse for a source detected in an OTELO\(_{\text{deep}}\) image. The classification presented in this work refers to the object closest in position (which usually also is the brightest object) detected in the $F814W$ image of HST-ACS. The detailed morphological analysis of OTELO sources up to $z = 2$ is the scope of a forthcoming paper (Nadolny et al. 2020).

### 4. Analysis and discussion

#### 4.1. Mass-metallicity relation

In this section we compare our results with the high-$z$ SDSS, GAMA, and VVDS-Deep and Ultra-Deep samples (see Sect. 2.2 for details on the sample selection). Finally, a comparison is made with the local SDSS sample in order to shed light on the possible evolution of the MZR at the low-mass end.

In Fig. 6 we show the local SDSS sample of galaxies up to $z \sim 0.1$ together with higher redshift samples up to $z \sim 0.33$ from GAMA and SDSS with their respective second-order
polynomial fits $[12 + \log(O/H) = a + bx + cx^2$ where $x = \log M_*].$ Both high redshift samples were selected to be complete in mass and luminosity at a given redshift. It is worth noting that the OTELO sample is three orders of magnitude deeper in stellar mass than the two surveys at similar redshift. This is due to the depth of the OTELO survey (line flux down to $5.6 \times 10^{-19}$ erg$^{-1}$ s$^{-1}$ cm$^{-2}$), see Bongiovanni et al. (2020) and its potential as a blind spectroscopic survey to pre-select ELSs according to their PS and not only as a result of the colour excess in the colour-magnitude diagram. In Fig. 6 we also show the selected sample from the VVDS-Ultra-Deep survey: after our selection process (see Sects. 2.2 and 3.2), we are left with 10 sources from a total of 1125 in the redshift range $0.2 < z < 0.42$ with $z_{\text{mean}} = 0.26$, which is much wider than our SFG sample redshift range ($0.36 < z < 0.42$). The highest redshift in the selected VVDS-Ultra-Deep sample is $z \sim 0.33$. The VVDS-Ultra-Deep survey data were processed in a consistent way with our data: (i) H$\alpha$ fluxes were corrected for stellar absorption (Eq. (4)) and for intrinsic extinction (in this case, using the Balmer decrement); (ii) we estimated the N2 index, SFR (assuming a Chabrier IMF) and sSFR (using VVDS stellar masses), and (iii) we removed sources with $N2 > -0.4$ as possible AGNs. We note that this sample spans the same range of metallicities, but covers a smaller range in stellar masses. We argue that this shows the capacity of the OTELO survey to recover the low-mass end of the observed galaxy population, even with its limited area of 56 arcmin$^2$ and despite the limited spectral range (230 Å wide, centred at 9175 Å).

In order to shed light on the possible evolution of the MZR, we compared our result with the local SDSS sample. The fit from Tremonti et al. (2004) is valid in the range of $8.5 < \log M_* < 11.5$, while our sample reaches masses as low as $\log M_* \sim 6.8$ (where 63% of the sample with $\log M_* < 9$), therefore we did not use it. In Fig. 6 we show our second-order polynomial fit to the local SDSS sample. From our SFG sample we identify sources with $[\text{N} \text{II}] \lambda 6583$ flux below the OTELO line flux limit (low-flux sources). Because for these low-flux sources our metallicity estimate has a large uncertainty, we provide two MZR fits: the first is obtained using sources with $[\text{N} \text{II}] \lambda 6583$ flux that is higher than the OTELO line flux limit (points and dot-dashed line in blue), and the second includes low-flux sources (points and dot-dashed line in gold). For the two fits we limited the $c$ parameter to be not greater than the $c$ parameter form fit of the local SDSS sample. Additionally, we show in this figure the MZR second-order polynomial fits to the local ($z < 0.1$) and high-$z$ ($z > 0.31$) SDSS and GAMA ($z \sim 0.32$) samples.

Our sample lies within the distribution of local SDSS MZR. Although our fit to sources whose $[\text{N} \text{II}] \lambda 6583$ line flux exceeds the OTELO line flux limit is flatter at the low-mass end, we cannot conclude that there is an evolution of the MZR because of the uncertainty in the metallicity of the low-flux sources. In the cases of the fit that include low-flux sources (golden symbols),

![Fig. 6. Mass-metallicity relation for OTELO H$\alpha$ ELS. Blue and golden markers represent sources with an $[\text{N} \text{II}] \lambda 6583$ line flux above and below OTELO flux limit, respectively. The light green empty square shows the MZR outlier as defined by Peeples et al. (2008). Cyan triangles show the VVDS-Ultra-Deep selected sample with $z_{\text{mean}} = 0.26$. Grey filled contours show local SDSS SFGs. Green filled contours show the high-$z$ SDSS sample. Red filled contours show the SFG sample from the GAMA survey. All fits shown in this figure represent second-order polynomial fits. The blue dot-dashed line is the OTELO MZR fit to the sources with an $[\text{N} \text{II}]$ line flux above the OTELO line flux limit, and the golden dot-dashed line shows the OTELO MZR including low-flux sources; the black thin dashed line shows the fit to the local SDSS sample; and green and red dashed lines show the SDSS and GAMA samples at $z \sim 0.3$, respectively. The error bars correspond to the 25% (red) and 68% (black) confidence interval around the most probable value. The calibration uncertainty in metallicity is not included.

Table 2. Polynomial coefficients of the MZR fits shown in Fig. 6 in the form $12 + \log(O/H) = a + bx + cx^2$, where $x = \log M_* .$

<table>
<thead>
<tr>
<th>Sample</th>
<th>$a$</th>
<th>$b$</th>
<th>$c$</th>
</tr>
</thead>
<tbody>
<tr>
<td>OTELO SF $z = 0.38$</td>
<td>6.103</td>
<td>0.494</td>
<td>-0.026</td>
</tr>
<tr>
<td>OTELO SF low-lim</td>
<td>4.836</td>
<td>0.620</td>
<td>-0.026</td>
</tr>
<tr>
<td>SDSS $z &lt; 0.1$</td>
<td>4.798</td>
<td>0.644</td>
<td>-0.026</td>
</tr>
<tr>
<td>SDSS $z = 0.31$</td>
<td>1.987</td>
<td>1.169</td>
<td>-0.051</td>
</tr>
<tr>
<td>GAMA $z = 0.32$ $^{(1)}$</td>
<td>-0.752</td>
<td>1.723</td>
<td>-0.079</td>
</tr>
</tbody>
</table>

Notes. $^{(1)}$Coefficients from second-order polynomial fit of the GAMA data below $z = 0.2$. 

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our fit approximates to the local SDSS MZR, especially at the low-mass end. All the fitted polynomial coefficients are listed in Table 2.

In Fig. 6 the upper histogram represents the mass distribution where the red dashed line shows the pure SFG sample and the grey filled histogram also includes possible AGNs (which are not included in the MZR). The “gap” in the \( M_* \) distribution between \( \sim 10^{8.2} \) and \( 10^9 M_\odot \) is a statistical effect due to the low number of sources in the sample. The right-hand histogram shows the metallicity distribution. The peak around \( z \approx 7.6 \) is associated with low-flux sources with an \([\text{NII}]\lambda 6583\) line flux that is lower than the OTELO flux limit. This is the hard limit of our method when it recovers line fluxes from PS, and also for the use of N2 as a metallicity estimate.

4.2. Mass, metallicity, and SFR

As proposed by Lara-López et al. (2010b), the stellar mass, metallicity, and SFR are correlated and may be reduced to a plane in 3D space. In this section we study the scaling relationships between SFR, sSFR, \( M_* \), and metallicity, shown in Fig. 7, where we represent the selected sub-samples from local SDSS, GAMA, VVDS (Deep and Ultra-Deep), and our \( \text{H}\alpha \) SFG sample. The left panel of this figure shows the dependence of the SFR on the stellar mass. Noeske et al. (2007) explored this relation over a wide range of redshifts between \( 0.2 < z < 1.1 \), with the mass complete down to \( \sim 10^{10} \). They found a main sequence for SFGs over cosmic time with a slope of 0.67. Our sample reaches three orders of magnitude lower in stellar mass than the work of Noeske et al. (2007). We find a linear relation with a slightly higher slope of 0.78 ± 0.11, but shifted down by \( -0.25 \) dex. This is associated with the different mass ranges in the two datasets. Using data from merged VVDS Deep and Ultra-Deep survey samples, we obtain a linear fit with a slope of 0.61 ± 0.03, shown in the same figure. We note that this fit approximates ours at the high-mass end. Our fit is in agreement with previous studies within the estimated dispersion. We find no evolution of this relation from the local Universe to \( z = 0.4 \). Our \( \text{H}\alpha \) sample is consistent with the local SDSS distribution, as can be seen not only in the spatial distribution, but also in the linear fits of the two samples. Our sample has moderate SFRs, ranging from \( -2.2 \leq \log \text{SFR} \leq 0.2 M_\odot \text{yr}^{-1} \). VVDS covers this relation with a similar dispersion from the high-mass (Deep) down to the low-mass (Ultra-Deep) regime; the VVDS-Ultra-Deep sample reaches the mean mass of the OTELO \( \text{H}\alpha \) sample (\( M_* \approx 10^8 M_\odot \)).

The middle and right panels of this figure show the relationship of sSFR with \( M_* \) and metallicity, respectively. There is no clear evidence of evolution as compared with the local SDSS sample. In both cases the sSFR is anti-correlated with \( M_* \) and metallicity (middle and right panel, respectively). Furthermore, these relations hold up to \( z = 0.4 \) down to the low-mass regime for the OTELO and VVDS samples.

4.3. Mass-size relation

In Fig. 8 we show the \( M_* - r_e \) relation (MSR) for the \( \text{H}\alpha \) sample we studied. We estimated the physical half-flux radius \( r_e \) using the effective radius from the Sérsic model and \( r_e \) values are given in the last column of Table B.1). From a linear fit we obtain an intercept of \( -1.41 \pm 0.73 \) and a slope of 0.19 ± 0.09. In this fit we used 13 sources from the SFG sample that had available \( M_* \) and \( r_e \) measurements. Here we compare our MSR with Ichikawa et al. (2012), who studied galaxies from the Moirirs Deep Survey (Kajisawa et al. 2009) in the GOODS-North region at \( 0.3 < z < 3 \) and in a mass range of \( M_* \sim 10^8 - 10^{11} M_\odot \). They found a universal slope of the MSR (independent of redshift and galaxy activity, SFGs or quiescent, for a given mass) of \( r_e \sim R_e^* \) with \( a \sim 0.1 - 0.2 \). Figure 8 shows a fit from Ichikawa et al. (2012) for colour-selected SFGs at \( 0.25 < z < 0.5 \) with a slope of 0.115, using their Eq. (3) with the parameters given in their Table 1. In Fig. 8 we also show a linear fit from van der Wel et al. (2014) for colour-selected SFGs from the 3D-HST/CANDELS at \( z \sim 0.25 \) (equation and values as given in their Table 1). We found that our results agree with those from Ichikawa et al. (2012). We argue that this is due to the mass range studied in both works.

4.4. Morphology

In Fig. 9 we show the results of our morphological classification using an automatic fitting of a single-Sérsic model and a visual classification. More than a half (12 sources, 57%; see the background histogram) of the full \( \text{H}\alpha \) sample are classified as discs with a median Sérsic index \( n \) of 1.31. Two (\( \sim 9.5 \% \)) objects are classified as spheroidal with \( n \) of 2.2. The increase of the Sérsic index from discs to spheroids is expected because early-type galaxies are fitted with higher Sérsic profile than late-type galaxies.
The percentage of metal-rich dwarf galaxies selected from the SDSS survey from Tremonti et al. (2004). They predicted that in order to have such high metallicity, these sources also need to have a lower gas fraction than objects with similar SFR and luminosities. They discussed other possible explanations for this behaviour: environment (interactions), effective yields, and finally the possibility that these sources are transitional dwarf galaxies at the end of their period of star formation. The outlier in our sample has a low stellar-mass log $M_\ast = 8.2$, high metallicity $12 + \log(O/H) = 8.6$, low SFR of $\log$ SFR $\sim -1 M_\odot$ yr$^{-1}$, and a high sSFR of log SF$\sim -9.2$. More information, in particular spectroscopic observations, is needed to disentangle its morphology and to shed light on its real nature.

5. Conclusions

We studied the MZR and scaling relations between stellar mass, SFR, and sSFR, and the morphology of the low-mass H$\alpha$ sample drawn from the OTELO survey at $z \sim 0.4$. We reach three orders of magnitude deeper in stellar mass ($\sim 10^{6.8} < M_\ast/M_\odot < 10^{10}$ with 57% of the sample with $M_\ast < 10^9 M_\odot$) than the SDSS and GAMA samples at similar redshift. The PS (data-product of OTELO survey) gives us the possibility to efficiently select and study the metallicity of these low-mass ELSs. About 95% of our ELS sample was selected from spectral features present in PS through a semi-automatic procedure detailed in Sect. 2. If the colour-excess technique used in classical NB surveys had been used instead, only half of this sample (i.e. necessarily the sources with larger observed EWs) would be recovered. For more details, see OTELO-I and Bongiovanni et al. (2020).

Because of the OTELO instrumental limits [as seen in the N2 index and/or $12 + \log(O/H)$] and the associated errors in the recovery of the line fluxes and the low number of objects in the sample, we cannot conclude about a possible MZR evolution that would be due to cosmic variance. Furthermore, our results are consistent with those obtained using the local SDSS sample (see Fig. 6).

We also explored the SFR as a function of $M_\ast$, which has previously been studied by Noeske et al. (2007), who found the so-called main sequence. We showed that this relation holds for the sample we studied, but with a systematic shift by $\sim 0.25$ dex in SFR, which is associated with the different stellar mass ranges used in Noeske et al. (2007) and in this work. Our data agree with the VVDS-Deep and Ultra-Deep data used in this work. Furthermore, the global distribution of the OTELO SFG sample lies within the local SDSS distribution.

From the automatic and visual morphological classification (see Sect. 4.4), we find that the majority of sources in the sample are discs (57%), followed by spheroids (9.5%) with median Sérsic indices of 1.31 and 2.2, respectively. Two sources are classified as point-like and tadpole (one in each class, 9.5% in total).
No classification is given for 24% of our sample because they are not detected in the high-resolution HST-ACS F814W image, have no data, or because the sources are just too faint to provide any reliable classification.

We obtained the $M_*$-size relation for our H$_{\alpha}$ sample. Our results agree with those of Ichikawa et al. (2012), but have a slightly higher slope (0.14 and 0.115, respectively). The difference in the results in this work and in van der Wel et al. (2014) is probably due to the different range of stellar-masses that is covered in both studies (we reach almost three orders of magnitude deeper in $M_*$ than van der Wel et al. (2014)).

We identified one candidate (~6% of our SFG sample) as a transitional dwarf galaxy with low mass $M_*$ = 8.6 $M_\odot$ and high metallicity $12 + \log(O/H) = 8.64$. Considering the same criteria, we find less than 1% of these sources in the local-SDSS sample. Morphologically, this source is classified as a compact (~0.7 kpc) disc galaxy ($n = 1.25$) with blue ($g - i$) colour. The study of MZR, SFR, sSFR, and its colour and morphology suggest that this source is indeed a transitional dwarf galaxy at the end of its star formation period, as discussed in previous works (e.g. Peeples et al. 2008). More information is needed to conclude on this issue, however.

Acknowledgements. This work was supported by the project Evolution of Galaxies, of reference AYA2014-88661-C3-1-P and AYA2018-88007-C3-2-P, within the “Programa estatal de fomento de la investigación científica y técnica de excelencia del Plan Estatal de Investigación Científica y Técnica de Innovación (2013-2016)” of the “Agencia Estatal de Investigación del Ministerio de Ciencia, Innovación y Universidades”, and co-financed by the FEDER “Fondo Europeo de Desarrollo Regional”. AG and MC are also supported by the Spanish State Research Agency grant MDM-2017-0737 (Unidad de Excelencia María de Maeztu CAB). JAD is grateful for the support from the UNAM-DGAPA-PASPA 2019 program, and the kind hospitality of the IAC. Based on observations made with the Gran Telescopio Canarias (GTC), installed in the Spanish Observatorio del Roque de los Muchachos of the Instituto de Astrofísica de Canarias, on the island of La Palma. MALL is a DArk-Carlsberg Foundation Fellow (Semper Ardens project CF15-0384). This research uses data from the VIMOS VLT Deep Survey, obtained from the VVDS database operated by Cesam, Laboratoire d’Astrophysique de Marseille, France. Based on observations obtained with MegaPrime/MegaCam, a joint project of CFHT and CEA/IRFU, at the Canada-France-Hawaii Telescope (CFHT) which is operated by the National Research Council of Canada, the Institut National des Sciences de l’Univers of the Centre National de la Recherche Scientifique (CNRS) of France, and the University of Hawaii. This work is based in part on data products produced at Terapix available at the Canadian Astronomy Data Centre as part of the Canada-France-Hawaii Telescope Legacy Survey, a collaborative project of NRC and CNRS. This work had made extensive use of TOPCAT and STILTS software (Taylor 2005).

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This work was supported by the project Evolution of Galaxies, of reference AYA2014-88661-C3-1-P and AYA2018-88007-C3-2-P, within the “Programma estatal de fomento de la investigación científica y técnica de excelencia del Plan Estatal de Investigación Científica y Técnica de Innovación (2013-2016)” of the “Agencia Estatal de Investigación del Ministerio de Ciencia, Innovación y Universidades”, and co-financed by the FEDER “Fondo Europeo de Desarrollo Regional”. AG and MC are also supported by the Spanish State Research Agency grant MDM-2017-0737 (Unidad de Excelencia María de Maeztu CAB). JAD is grateful for the support from the UNAM-DGAPA-PASPA 2019 program, and the kind hospitality of the IAC. Based on observations made with the Gran Telescopio Canarias (GTC), installed in the Spanish Observatorio del Roque de los Muchachos of the Instituto de Astrofísica de Canarias, on the island of La Palma. MALL is a DArk-Carlsberg Foundation Fellow (Semper Ardens project CF15-0384). This research uses data from the VIMOS VLT Deep Survey, obtained from the VVDS database operated by Cesam, Laboratoire d’Astrophysique de Marseille, France. Based on observations obtained with MegaPrime/MegaCam, a joint project of CFHT and CEA/IRFU, at the Canada-France-Hawaii Telescope (CFHT) which is operated by the National Research Council (NRC) of Canada, the Institut National des Sciences de l’Univers of the Centre National de la Recherche Scientifique (CNRS) of France, and the University of Hawaii. This work is based in part on data products produced at Terapix available at the Canadian Astronomy Data Centre as part of the Canada-France-Hawaii Telescope Legacy Survey, a collaborative project of NRC and CNRS. This work had made extensive use of TOPCAT and STILTS software (Taylor 2005).
Appendix A: Inverse deconvolution and uncertainty estimates

Fig. A.1. Observed PS, the best-fit PS, and the envelope of the PS simulations where all parameters are in the 25% confidence interval (light green) and 68% confidence interval (light grey) for the first eight SFG for which an $H\alpha + [N\text{II}]$ line is detected. Violet, blue, and cyan vertical lines in the PS figures show the position of the $[N\text{II}]\lambda 6548$, $H\alpha$, and $[N\text{II}]\lambda 6583$ line, respectively, for the best $z$ obtained by the fit. The red points that appear in some observed PS were removed during the fitting process because they correspond to the $[S\text{II}]\lambda\lambda 6716,6731$ doublet or may be artefacts. To the right of each PS is the PDF of $12 + \log(O/H)$ that was obtained from the Monte Carlo simulations where the mode (our reference value), the 25%, 68%, and 90% confidence intervals around the mode are shown as vertical lines; the black vertical line in these plots shows the empirical division between SFR and AGN at $12 + \log(O/H) = 8.67$ ($N2 = -0.4$).

Throughout this paper we made use of the inverse deconvolution of the pseudo-spectra (PS). In this appendix we describe how this deconvolution was performed.
the deviation $\sigma_{\text{med,c}}$ around this value (mean of the square root of the differences around $f_i$). We also have a redshift guess value $z_{\text{GUESS}}$ that is obtained from visual inspection of the PS in our GUI web-based tool, and it is assumed to be accurate up to the third digit, that is, with an uncertainty $\Delta z = 0.001$. Finally, we define the wavelength region of the PS where the lines can be detected in such a way that it contains at least 2.5 slices (15 Å being the average width of the RFT) below the redshifted [N II] 6548 line and 2.5 slices above the redshifted [N II] 6583 line:

$$\lambda_{\text{min}} = 6548 \cdot (1 + z_{\text{GUESS}} - \Delta z) - 15$$

$$\lambda_{\text{max}} = 6583 \cdot (1 + z_{\text{GUESS}} + \Delta z) + 15.$$  

(A.1)

We assumed a model spectrum as a rest-frame spectrum defined by Gaussian profiles of the [N II] 6548, H$\alpha$, and [N II] 6583 lines defined by their amplitude and a common line width for three lines, and a constant continuum level. The model to be compared with the observational data is therefore defined by following parameters: $f_{\text{mod}}, f_\alpha, \sigma_{\text{mod,c}}, f_{\text{PS}}^{\text{max}}$, and $\sigma_{\text{PS}}^{\text{max}}$ to constrain the parameter space (see below). Each simulation was convolved with the RFT filter system, and a synthetic PS $f_{\text{PS,mod},i}$ for each of the $N$ slices of the PS was obtained. This synthetic PS was compared with the observed PS, and the following likelihood function was obtained for each model and was compared with each point $i$ in the PS:

$$L(f_{\text{mod}}|f_{\text{PS}}) \propto \Pi_{i=1}^N w_i \exp \left( -\frac{\chi_i^2}{2} \right),$$

where $\chi^2$ is given in the usual way,

$$\chi_i^2 = \left( \frac{f_{\text{mod},i} - f_{\text{obs},i}}{\sigma_{\text{obs},i}} \right)^2,$$

(A.3)

and $w_i$ is a normalized weight function defined as

$$w_i \begin{cases} \text{MAX} \left( 1, \frac{f_{\text{mod,i}}}{f_{\text{obs,i}}} \right) & \text{if } \lambda_i \in [\lambda_{\text{min}}, \lambda_{\text{max}}] \\ 1 & \text{otherwise} \end{cases}.$$  

(A.4)

The use of the weight function aims to give a greater weight in the fit to the points where the line is presented and are above the continuum level, and in particular, to fit the core of the H$\alpha$ line. In practical cases, the weight function only has an effect in the objects with faint lines.

In order to sample the likelihood function, the first 30% (300 000) of the simulations made use of the following priors:

1. The redshift $z$ follows a flat distribution in the range $z_{\text{GUESS}} \pm 0.001$.
2. The dispersion of the line, $\sigma_{\text{line}}$, varies following a flat distribution from 20 to 500 km s$^{-1}$, which is translated into a range of 0.43–11 Å for the H$\alpha$ wavelength. This value is used for the three lines.
3. The continuum value follows a flat distribution in the interval [MAX(0, $f_{\alpha} - \sigma_{\text{med,c}}$), $f_{\alpha} + \sigma_{\text{med,c}}$].
4. The rest-frame $f_{\text{H$\alpha$}}$ line flux varies following a flat distribution in the interval [$f_{\alpha}, f_{\text{max}}$], where $f_{\text{max}}$ is given by

$$f_{\text{max}} = (f_{\text{mod,c}} + \sigma_{\text{PS}}^{\text{max}} + 3 \sigma_{\text{PS}}^{\text{max}}) \cdot (1 + z_{\text{max}}) \cdot \sqrt{2\pi}.$$  

(A.5)

5. The rest-frame $f_{\text{N II,6583}}$ line flux varies following a flat distribution between 0.006 $f_{\text{H$\alpha$}}$ and $f_{\text{max}}$.
6. The $f_{\text{N II,6548}}$ line flux is fixed to 1/3 $f_{\text{H$\alpha$}} f_{\text{N II,6583}}$.

When these 300 000 simulations were obtained, we computed a preliminary likelihood function. This likelihood was sampled in a regular interval with 200 bins over the range in which each parameter varied.

This likelihood function was marginalized for each parameter (i.e. we made a projection of the PDF on each particular parameter and the corresponding PDF was normalized). We obtained the marginalized likelihood using a regular interval of 200 bins over the range in which each parameter varied, and we obtained the range that contains 99% of the simulations around the value with the maximum likelihood. The next 20% (200 000) of the simulations were restricted to the parameters that sample this 99% confidence interval. This process was repeated five more times (using 10% (100 000) of the simulations at each step) with the limits obtained by the 85%, 75%, 68%, 40%, and 10% confidence intervals around the most probable value.

After this process we had a final PDF shape for each parameter, as well as all the required additional quantities such as the EWs, the N2 parameter, or the metallicities. We note that the asymmetry of the PDF implies that the mean and the standard deviation (squared of the variance) are not informative to describe the distribution. For this reason we used the mode of the resulting PDF as the reference value, and we obtained the region around the mode that included 10%, 25%, 50%, 68%, and 90% of the area. In the main text the error bars show the 25% and 68% confidence interval, which also gives an idea of the asymmetry of the associated distribution.

We show in Figs. A.1 and A.2 the observed PS, the best-fit PS, and the envelope of the PS simulations where all parameters are in the 25% confidence interval (light green) and 68% confidence interval (light grey) for the SPG for which the H$\alpha$ + [N II] line was detected. We also show the PDF of 12 + log(O/H) that we obtained from the Monte Carlo simulations for each source where the mode (our reference value), the 25%, 68%, and 90% confidence intervals around the mode are shown in vertical lines. We also show with a black vertical line the empirical division between SFG and AGN at 12 + log(O/H) = 8.67 (N2 = −0.4). Figures A.3 and A.5 show similar information for the possible transitional dwarf galaxy id: 3137, and for sources whose [N II] line flux is lower than the OTELO line flux limit (low-flux sources, respectively); we note that for the latter, the PDF of 12 + log(O/H) has a step-like shape that reflects the large uncertainty in the [N II] line flux. Finally, Fig. A.4 show the same information for AGN sources; we kept the PDF of 12 + log(O/H) for comparison with the previous plots.

We show in Table B.1 the modal values and 68% confidence intervals around the mode for $f_{\text{H$\alpha$}}$ and EW(H$\alpha$) (after correction for stellar absorption), $f_{\text{N II,6583}}$, and 12 + log(O/H). See Sect. 3.2 for details about the uncertainties of $z_{\text{OTELO}}$.  

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J. Nadolny et al.: MZR of the OTELO H$\alpha$ sources
Fig. A.2. As Fig. A.1 for the remaining two sources with a detected [N II] line.

Fig. A.3. As Fig. A.1, but for the possible transitional dwarf galaxy (source id: 3137).

Fig. A.4. As Fig. A.1, but for the three AGN candidates.
Appendix B: Uncertainty in stellar masses and SFR

As explained in Sect. 3, stellar masses were computed using the prescription of LS18, which makes use of colour ($g - i$) and the $i$ absolute magnitude. The associated uncertainty was computed making use of the R subroutine propagate, which makes an estimate of the error by a first- and second-order propagation of any input recipe, as well as by Monte Carlo simulations. The uncertainty quoted in Table B.1 corresponds to the variance obtained from 10 000 the Monte Carlo simulations using that sub-routine.

Finally, the uncertainty in the SFR quoted in Table B.1 requires some cautions. First of all, the SFR requires the use of an extinction value. We used the value that was obtained with the LePhare software and used empirical galaxy templates. This choice implies that the extinction value is a lower limit because the intrinsic extinction of the templates are unknown. In addition, LePhare does not provide an evaluation of the uncertainty in $E(B-V)$, therefore this uncertainty was not included in the SFR. Finally, we recall that our $f_{\text{H}\alpha}$ uncertainty does not correspond to a variance (where standard propagation theory would apply), but to a confidence interval. Taking these considerations into account, we computed the uncertainty in the SFR by applying the same formulae as we used to obtain the nominal SFR (the extinction correction used by Ly et al. 2012 and the SFR estimate using the Kennicutt & Evans 2012 relation) to the $f_{\text{H}\alpha}$ values that define the 68% confidence interval. We stress that the uncertainty values would be underestimated because they do not include all the possible sources of uncertainty (interstellar extinction in particular).
Table B.1. Catalogue of the OTELO Hα SFG sample. AGN candidates are separated in the lower part of the table.

<table>
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<tr>
<th>idobj</th>
<th>$E(B-V)$</th>
<th>$z_{\text{OTELO}}$</th>
<th>Mass</th>
<th>$12 + \log(O/H)$</th>
<th>$f_{\lambda_{6583}}$</th>
<th>$f_{\lambda_{6583}}$</th>
<th>EW(Hα)</th>
<th>SFR$_{\text{SFG}}$</th>
<th>$n$</th>
<th>$r_e$</th>
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<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>[10$^{-17}$ erg s$^{-1}$ cm$^{-2}$]</td>
<td>[10$^{-17}$ erg s$^{-1}$ cm$^{-2}$]</td>
<td>[Å]</td>
<td>[M$_{\odot}$ yr$^{-1}$]</td>
<td></td>
<td>[kpc]</td>
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Sources with [N II]6583 above OTELO line flux limit

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<td>625</td>
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<td>0.371</td>
<td>9.80</td>
<td>0.10</td>
<td>8.56±0.03</td>
<td>0.04</td>
<td>23.5±1.3</td>
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<td>1.12±0.06</td>
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<td>754</td>
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<td>0.12</td>
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<td>0.19±0.00</td>
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<td>0.39±0.13</td>
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<td>0.09±0.01</td>
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<td>8.16±0.07</td>
<td>0.10</td>
<td>32.0±0.31</td>
<td>1.60±0.34</td>
<td>88.2±1.8</td>
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<td>0.04</td>
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<td>8.43±0.17</td>
<td>0.11</td>
<td>0.41±0.06</td>
<td>0.06±0.06</td>
<td>126±0.30</td>
<td>0.01±0.04</td>
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<td>3106</td>
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<td>0.368</td>
<td>8.03</td>
<td>0.17</td>
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<td>0.05</td>
<td>0.52±0.10</td>
<td>0.07±0.04</td>
<td>48.5±1.6</td>
<td>0.01±0.03</td>
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<tr>
<td>3137</td>
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<td>0.389</td>
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<td>1.88±0.26</td>
<td>0.67±0.17</td>
<td>47.8±1.4</td>
<td>0.10±0.01</td>
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<td>7990</td>
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<td>0.11</td>
<td>8.30±0.16</td>
<td>0.16</td>
<td>0.63±0.08</td>
<td>0.05±0.04</td>
<td>35.1±0.7</td>
<td>0.18±0.02</td>
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<td>8074</td>
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<td>27.0±1.1</td>
<td>5.58±0.41</td>
<td>67.2±2.7</td>
<td>1.29±0.05</td>
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<td>10512</td>
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<td>0.380</td>
<td>8.05</td>
<td>0.11</td>
<td>7.92±0.10</td>
<td>0.01</td>
<td>3.29±0.12</td>
<td>0.06±0.06</td>
<td>231±10</td>
<td>0.094±0.003</td>
</tr>
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Sources whose [N II]6583 is lower than the OTELO line flux limit

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<td>3854</td>
<td>0.3</td>
<td>0.372</td>
<td>7.71</td>
<td>0.13</td>
<td>8.21±0.14</td>
<td>0.18</td>
<td>0.33±0.085</td>
<td>0.02±0.003</td>
<td>69±20</td>
<td>0.05±0.01</td>
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<tr>
<td>3855</td>
<td>0.0</td>
<td>0.373</td>
<td>7.35</td>
<td>0.12</td>
<td>7.62±0.57</td>
<td>0.05</td>
<td>0.22±0.07</td>
<td>0.01±0.041</td>
<td>37±16</td>
<td>0.006±0.002</td>
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<td>0.11</td>
<td>7.61±0.02</td>
<td>0.02</td>
<td>0.89±0.13</td>
<td>0.01±0.01</td>
<td>19.5±0.6</td>
<td>0.046±0.007</td>
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<tr>
<td>8713</td>
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<td>7.47</td>
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<td>0.2±0.05</td>
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<td>9932</td>
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<td>0.380</td>
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<td>1.40±0.12</td>
<td>0.008±0.008</td>
<td>71±7</td>
<td>0.04±0.003</td>
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AGN candidates

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<td>0.21</td>
<td>8.70±0.02</td>
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<td>0.61±0.06</td>
<td>0.20±0.05</td>
<td>163±16</td>
<td>0.032±0.003</td>
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</table>

Notes. Values (i.e. metallicity and SFR) for these sources need to be taken with caution. (1) Object selected via DEEP2 redshift, see Sect. 2.1 for details. (2) Possible transitional dwarf galaxy. (3) Uncertainty on $z_{\text{OTELO}}$ is ±0.001 (1 + $z_{\text{OTELO}}$). (4) Uncertainty defined as the 68% confidence interval around the mode. (5) Uncertainty in metallicity does not include calibration uncertainty. (6) Corrected for stellar absorption. (7) Lower limit; see Appendix B for the uncertainty evaluation.